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2

Single Chain Properties

2.1

The Ideal Polymer Chain

What Do They Have in Common?

a polymer chain



a drunk person



Brownian motion
(gas particles)

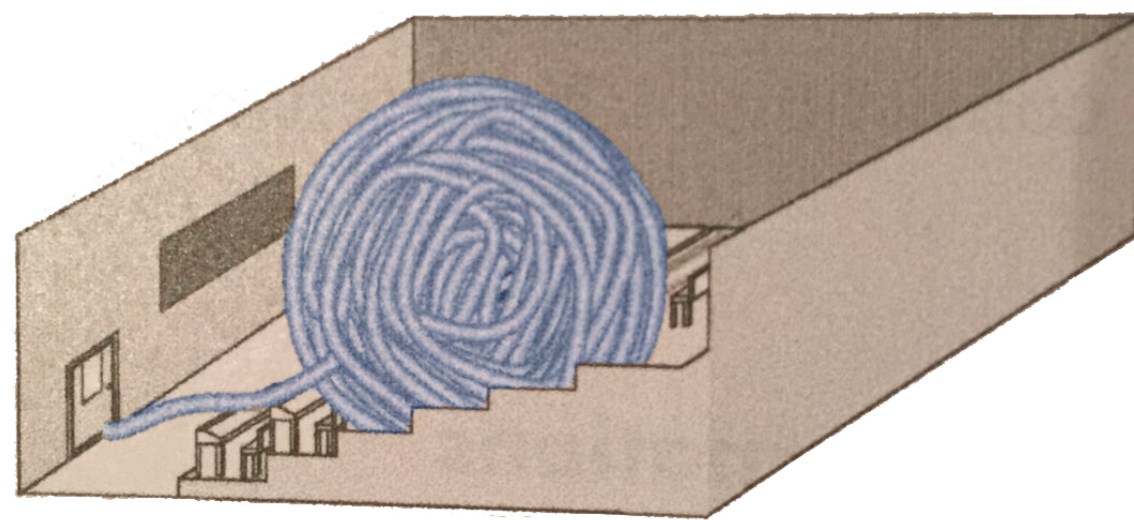


- ... can be described by the same statistical approach (the random walk model)
- How can this be used to express the polymer chain size? How does it relate to polymer molecular weight?

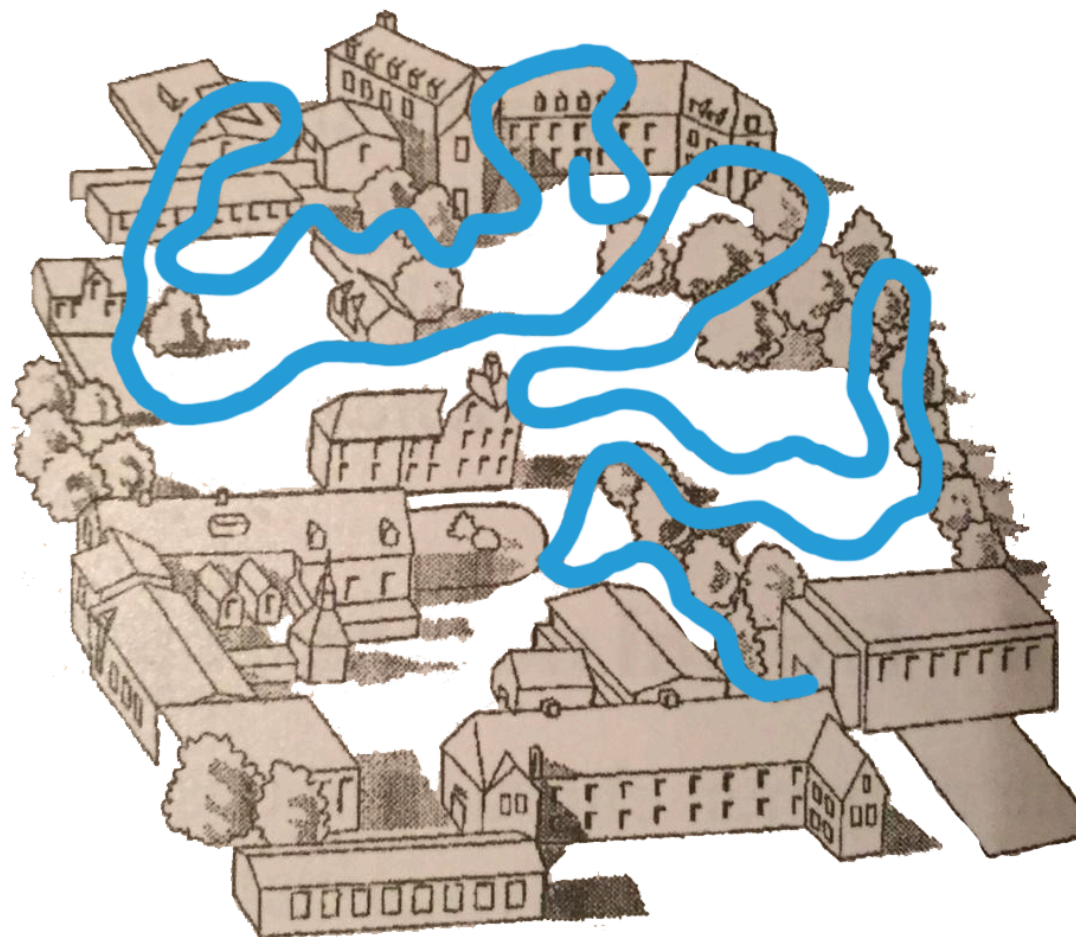
Chain Dimensions

- polymer conformation depends on chain flexibility, interactions between repeat units, and interactions with the surrounding medium
- example of a hypothetical polymer size: 10^{10} monomers; magnification factor: 10^8

collapsed globule
(attractive interactions)



random walk
(no effective interaction)



extended conformation
(long-range repulsion)

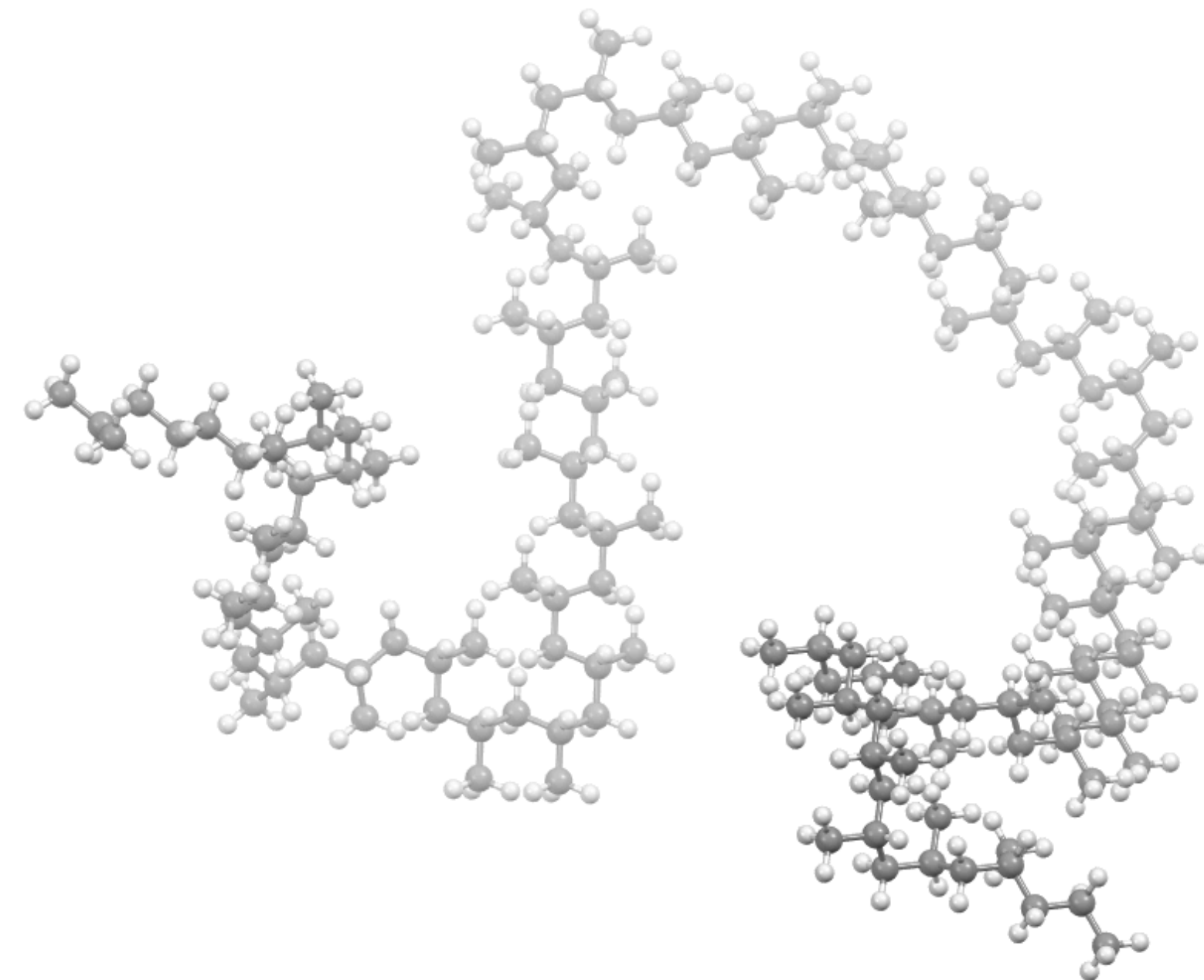


the ideal polymer chain

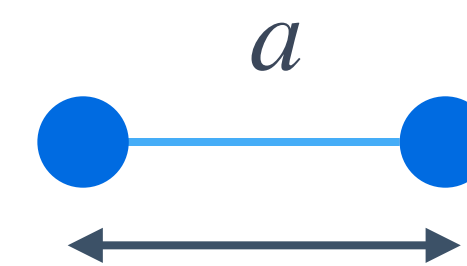
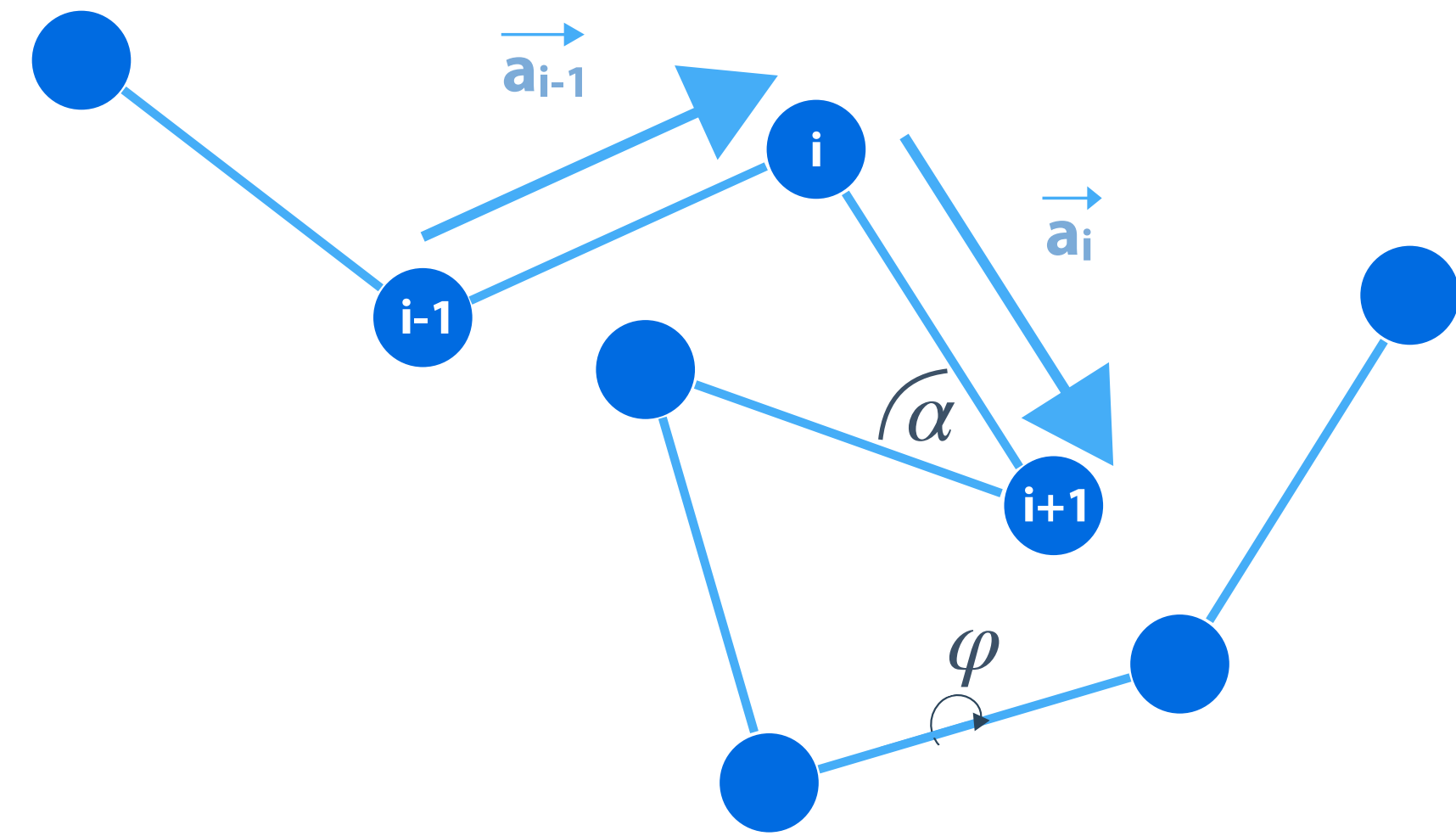
- chain sizes may be of vastly different dimensions, strongly depending on parameters like T or solvent

Simplification of the Polymer Chain

- from atoms and bonds to a chain of beads and links: models with variable repeating unit size, link sizes, bending and torsion angles



snapshot of PP (50 repeating units)



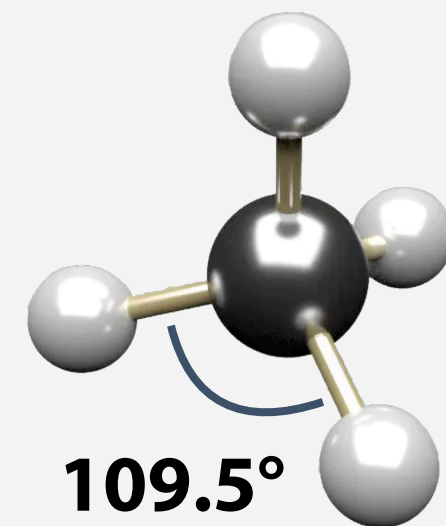
$$V_{stretch} = \frac{\kappa}{2}(a - a_0)^2$$

- ideal polymer chain: no energetic interactions between repeating units
- accurate description of polymer melts and solutions

Polymer Conformations

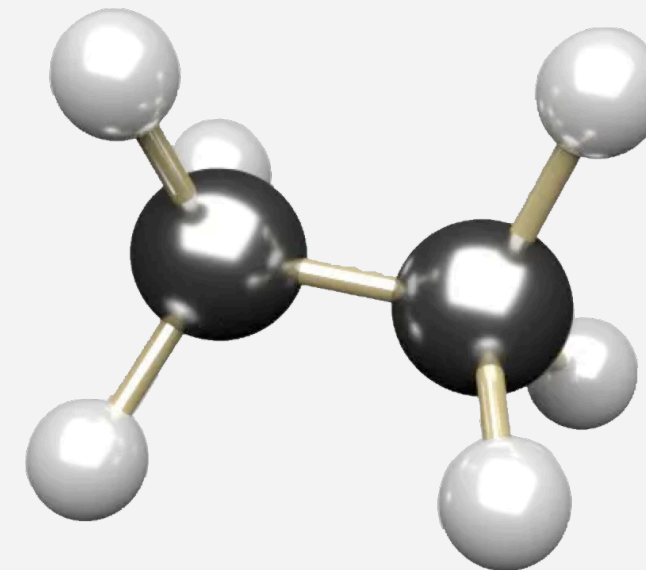
- conformation: the shape of a molecule resulting from rotation around fixed bonds

methane



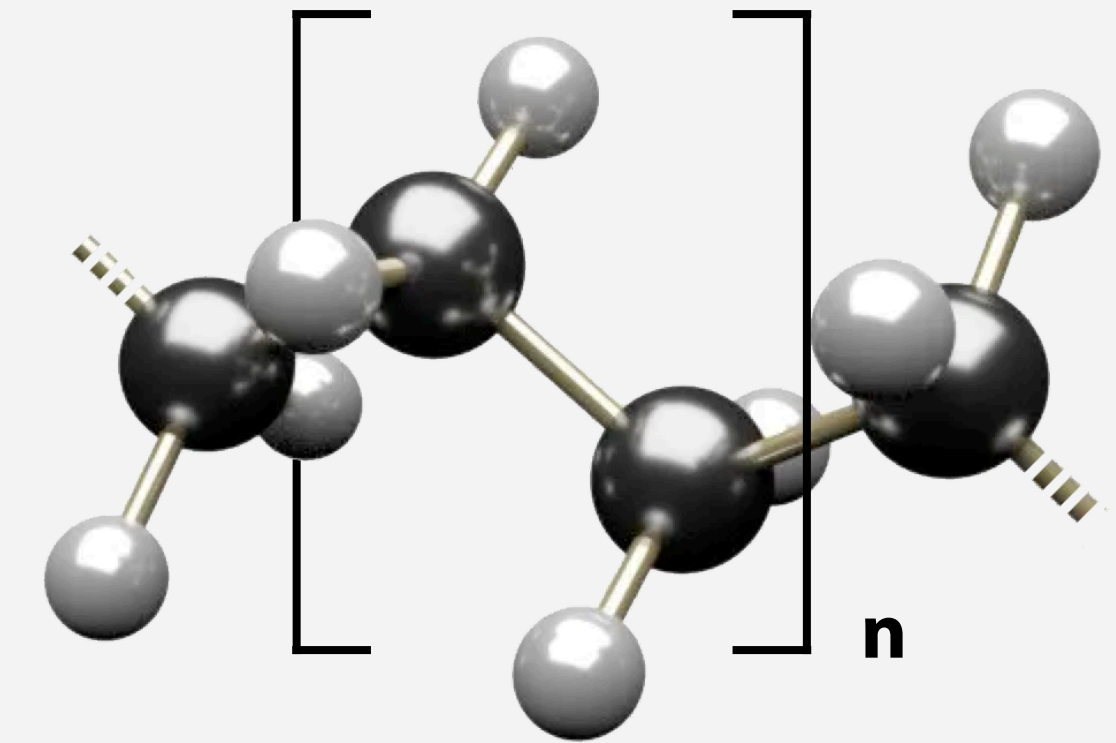
1 conformation
(fixed bond length
and angles)

ethane



3 conformations
(almost free rotation
around C-C bond)

polyethylene

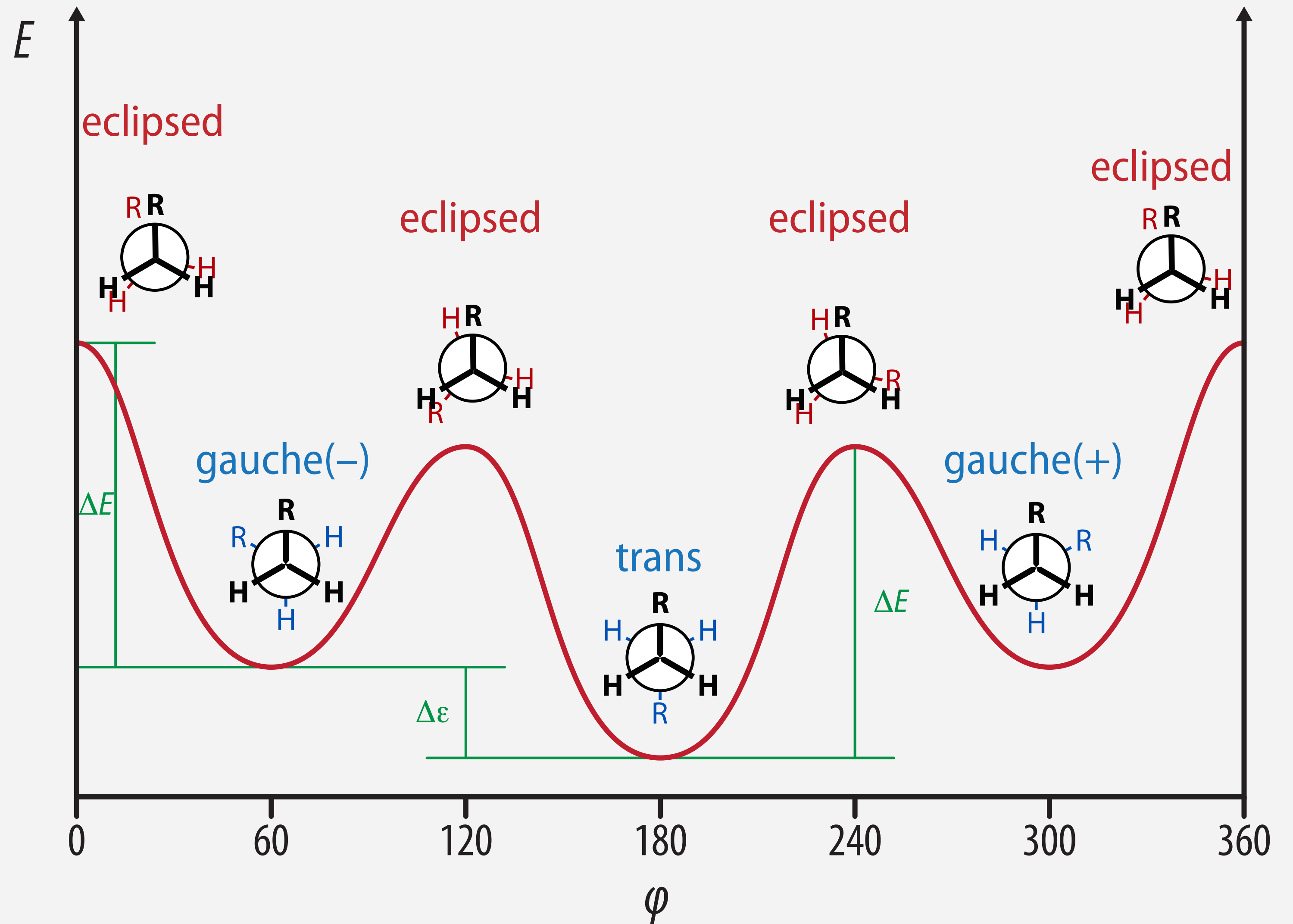


3^{2n} conformations

- the number of possible conformations increases drastically with the chain length
- conformational changes happen on the picosecond time scale

Energy Landscape of Polymer Conformations

main source of flexibility is the variation of torsion angles.



- rotation around single bonds at r.t. is governed by ΔE (thermal equilibrium) and $\Delta \epsilon$ (energy barrier)

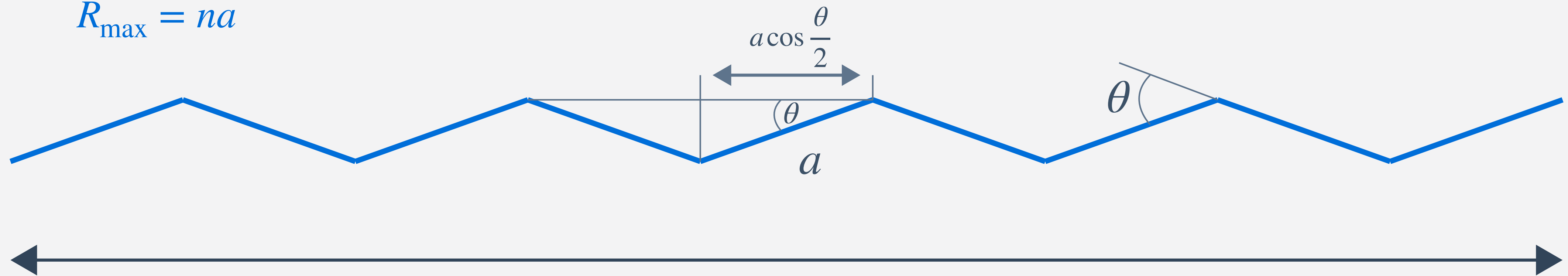
Contour Length

- the largest possible end-to-end distance (contour length or projection length), R_{max} , is an all-*trans* (zig-zag) conformation.

contour length:

$$R_{max} = na$$

all-*trans* conformation



projection length: $R_{max} = na \cos \frac{\theta}{2}$
(maximum end-to-end distance)

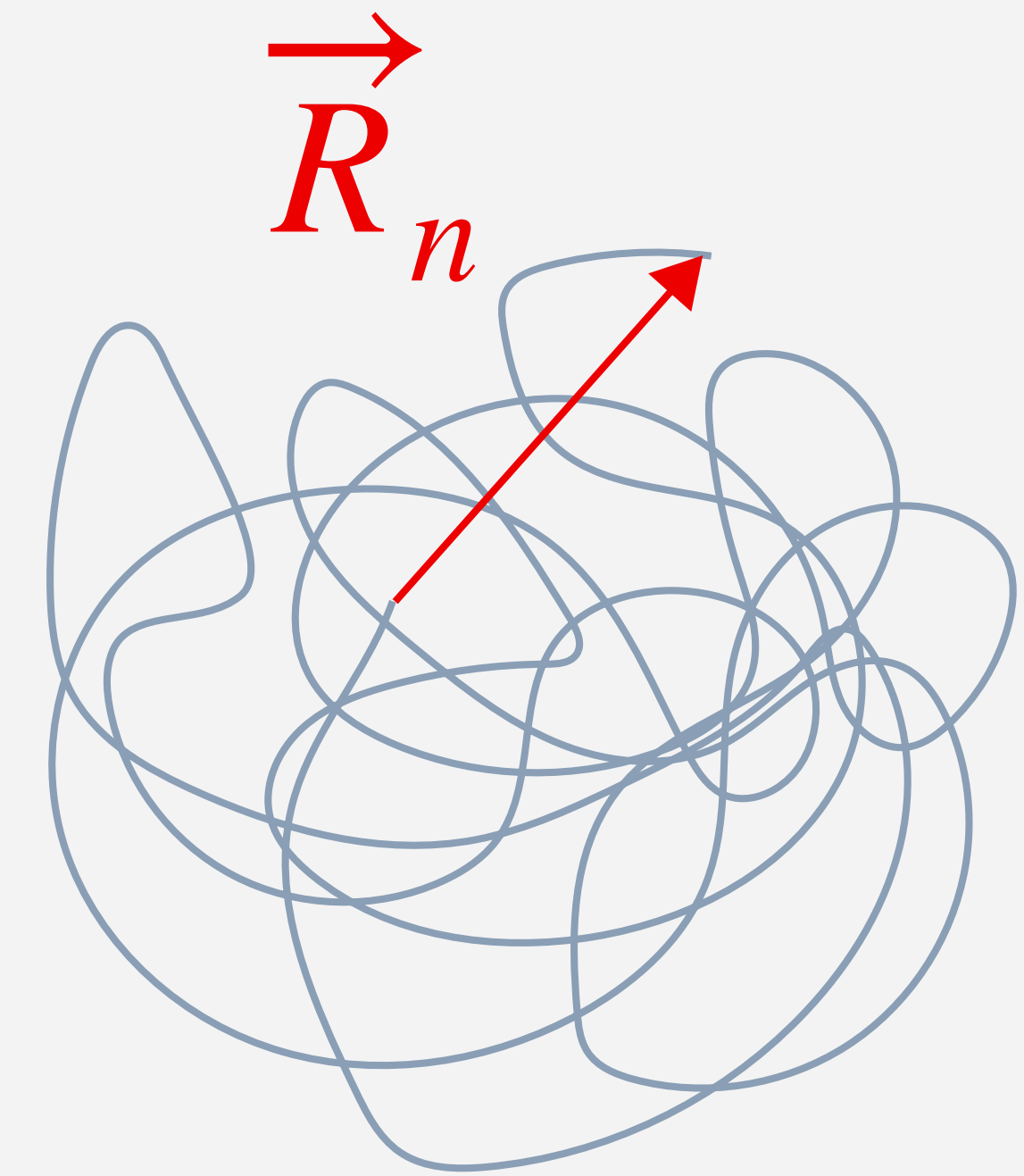
Example: PE

tetrahedral angle: $\theta = 68^\circ$,
bond length: $a = 1.54 \text{ \AA}$

- However, *gauche* states of torsion angles lead to flexibility in the chain conformation

Gauche and Trans States in PE

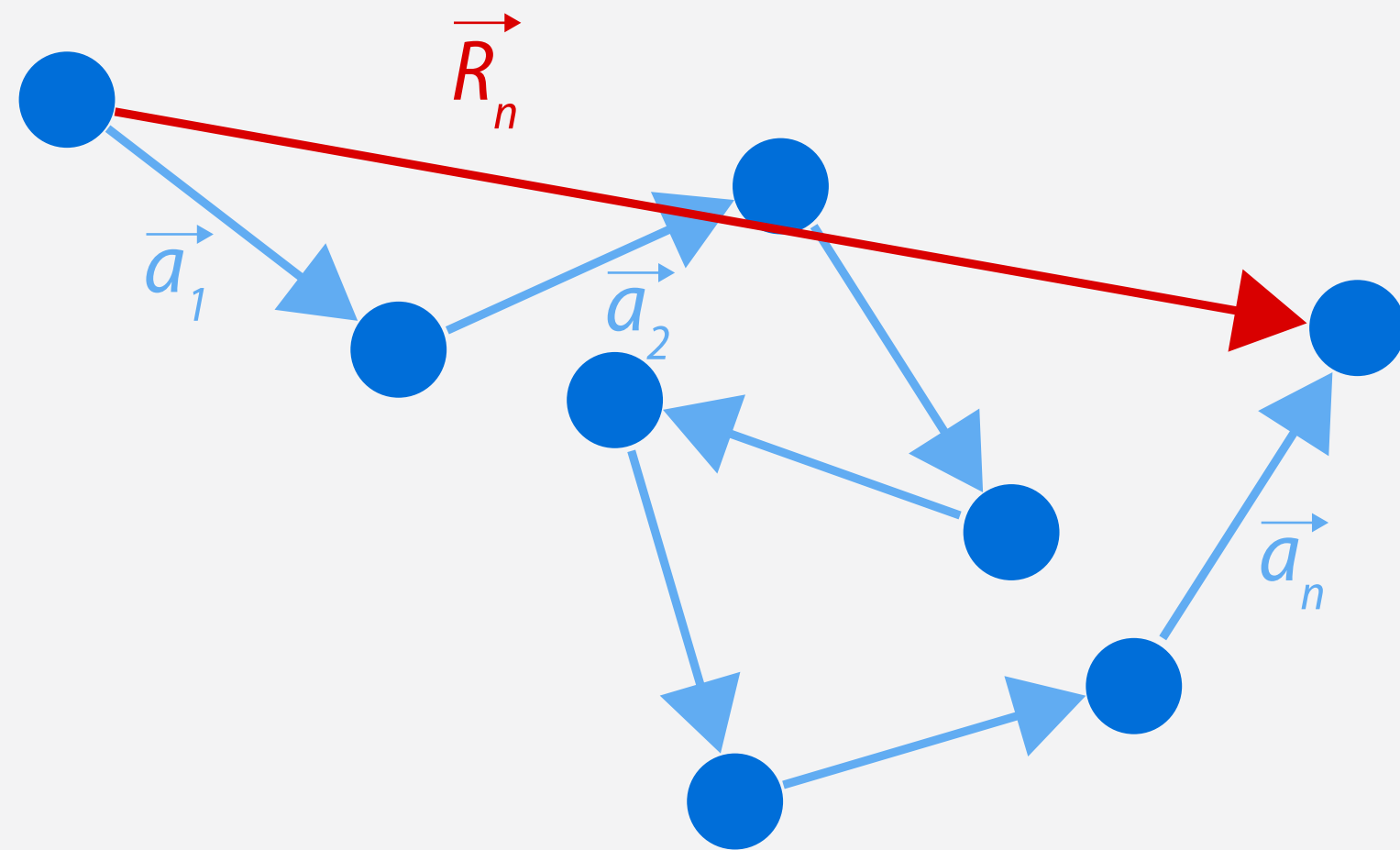
- polyethylene chain with 10'000 carbon atoms



- typically, all-*trans* rod-like chain sections comprise fewer than 10 main-chain bonds
- most synthetic polymers are hence quite flexible and are represented as random coils

Freely Jointed Chain Model

- no restriction upon bond angle and bond rotation
- however, $\langle \vec{R}_n \rangle$ is zero for an isotropic collection of ideal chains



$$\vec{R}_n = \vec{a}_1 + \vec{a}_2 + \dots + \vec{a}_n$$

polymer chain size

$$\langle R_n^2 \rangle^{\frac{1}{2}} \cong \sqrt{n} a$$

- polymer chain size represented by the mean-square end-to-end distance

$$\langle R_n^2 \rangle \equiv \langle \vec{R}_n \cdot \vec{R}_n \rangle = \left\langle \sum_{i=1}^n \vec{a}_i \cdot \sum_{j=1}^n \vec{a}_j \right\rangle = \sum_{i=1}^n a_i^2 + \sum_{i \neq j} \langle \vec{a}_i \cdot \vec{a}_j \rangle = n|a|^2 + \langle \sum_{i \neq j} \vec{a}_i \cdot \vec{a}_j \rangle = n|a|^2 \equiv na^2$$

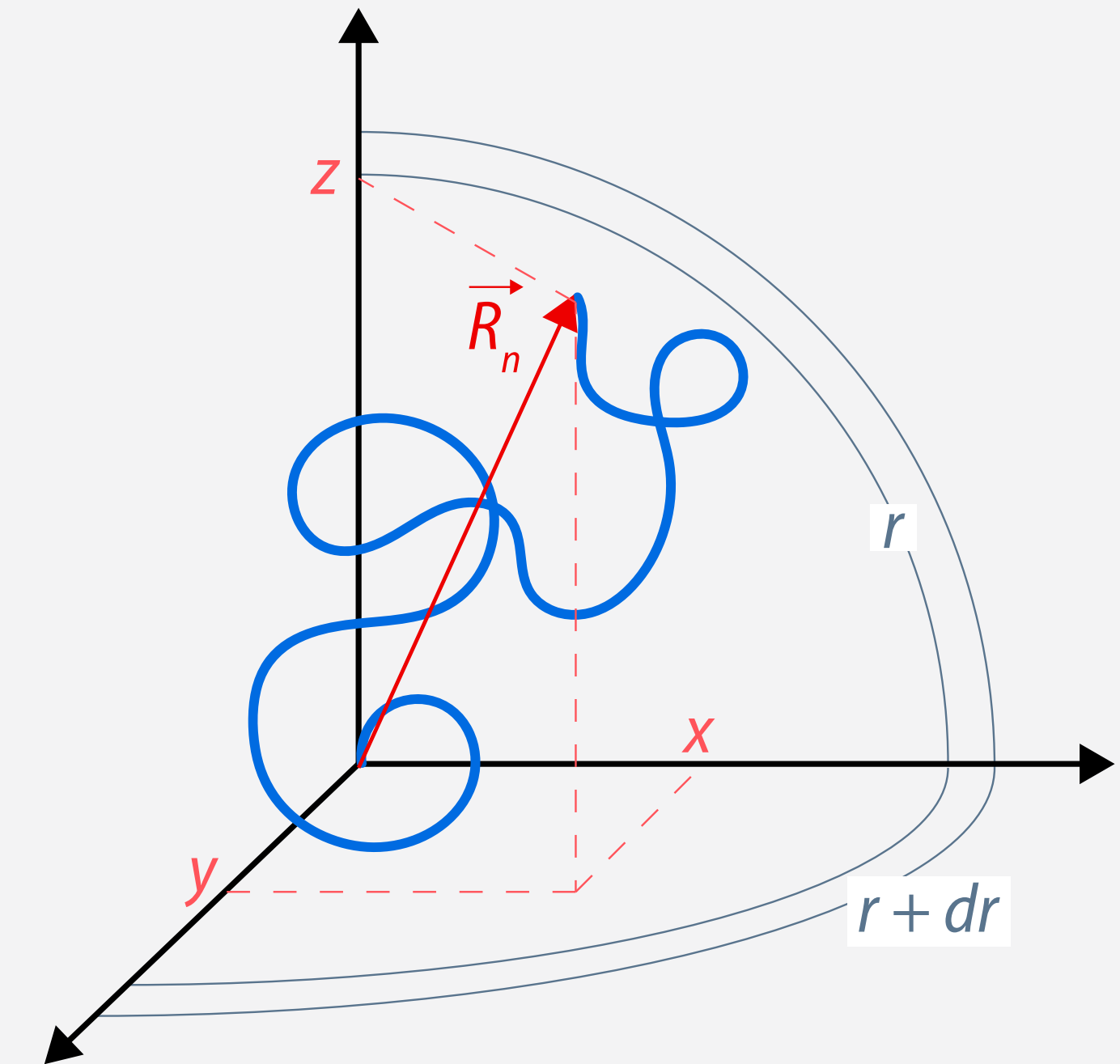
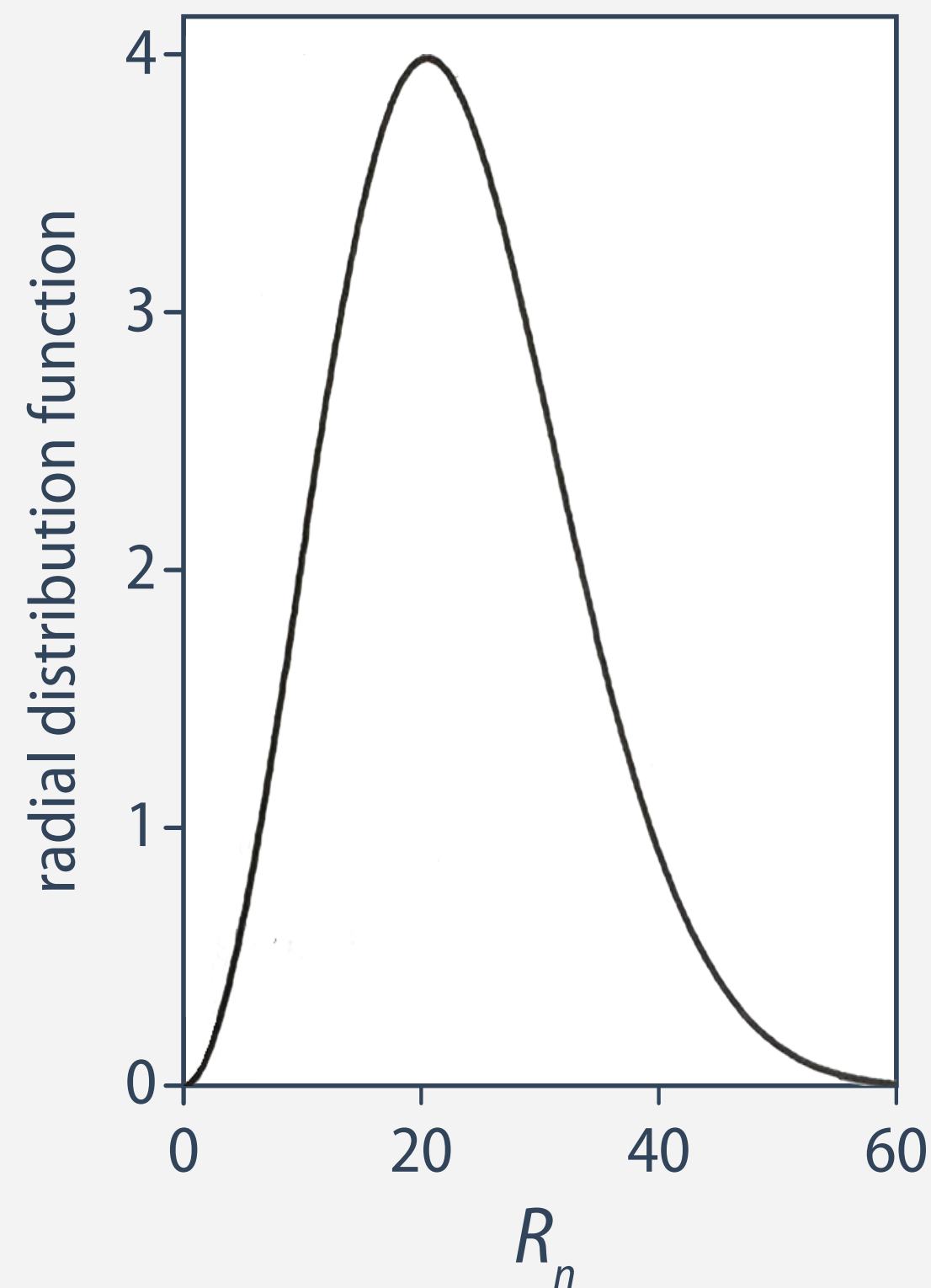
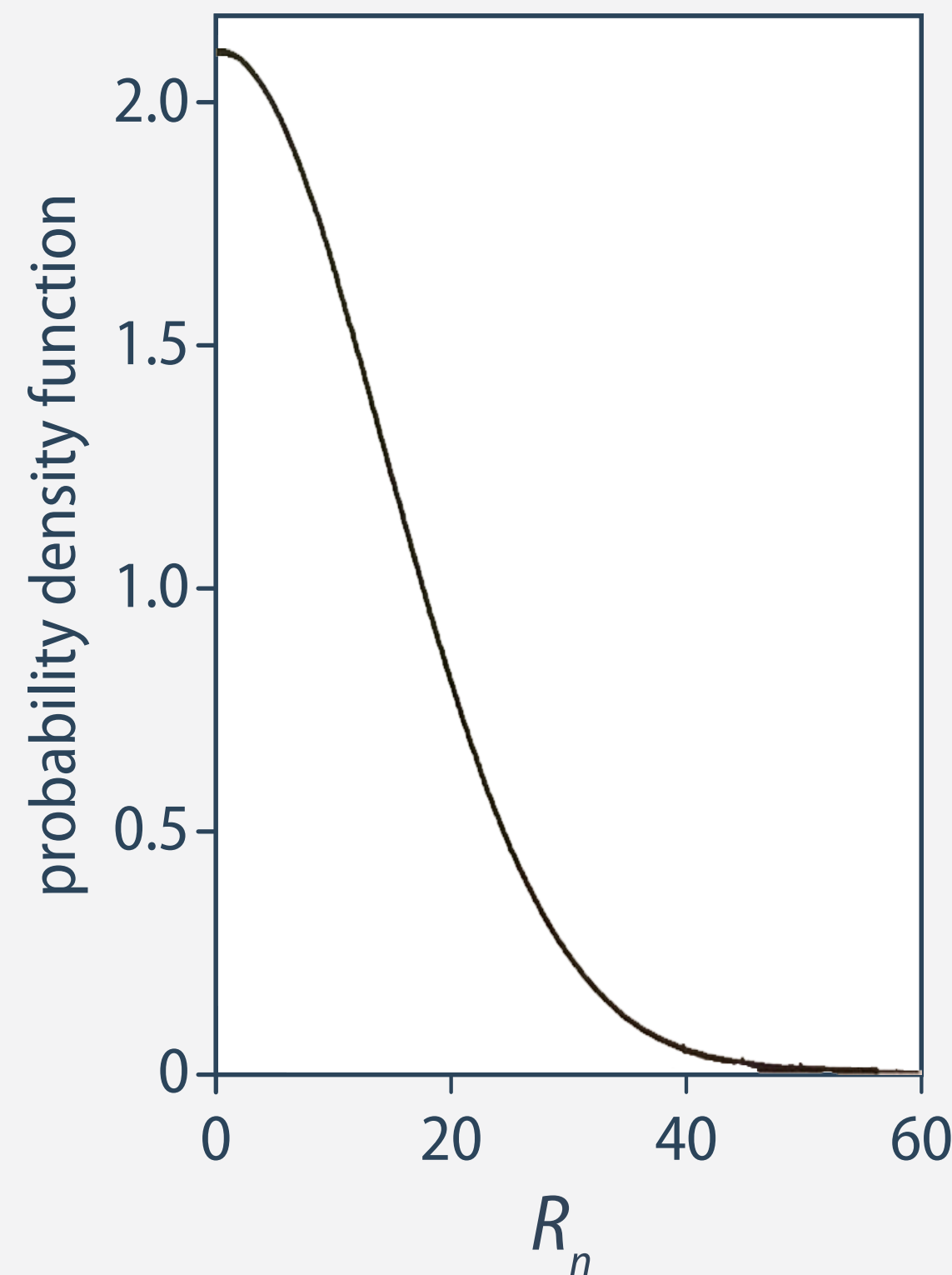
- hence, the root mean-square end-to-end distance is proportional to \sqrt{M} $\langle R^2 \rangle^{\frac{1}{2}} \propto \sqrt{n} \propto \sqrt{M}$

The “Gaussian” Chain

- an ideal chain can be mapped onto a random walk and obeys Gaussian statistics

$$P = \left(\frac{3}{2\pi na^2}\right)^{3/2} \exp\left(-\frac{3R_n^2}{2na^2}\right)$$

$$P4\pi R^2 dr = 4\pi \left(\frac{3}{2\pi na^2}\right)^{3/2} \exp\left(-\frac{3R_n^2}{2na^2}\right) R_n^2 dr$$



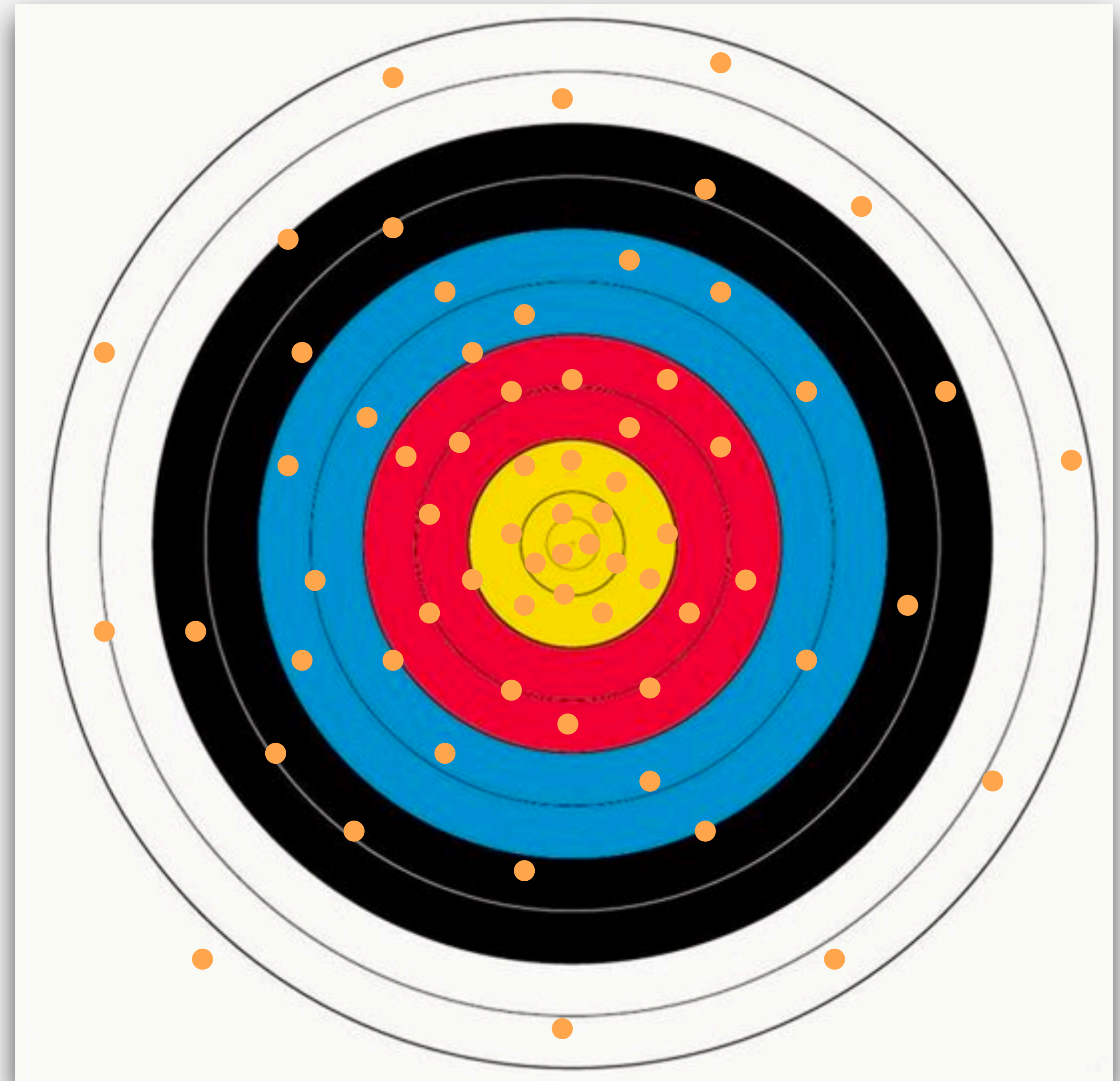
- most probable are conformations with $\vec{R}_n = 0$, but it's *rms* value is finite and proportional to \sqrt{n}

Analogy for Probability Density and Radial Distribution Function

- consider archery:

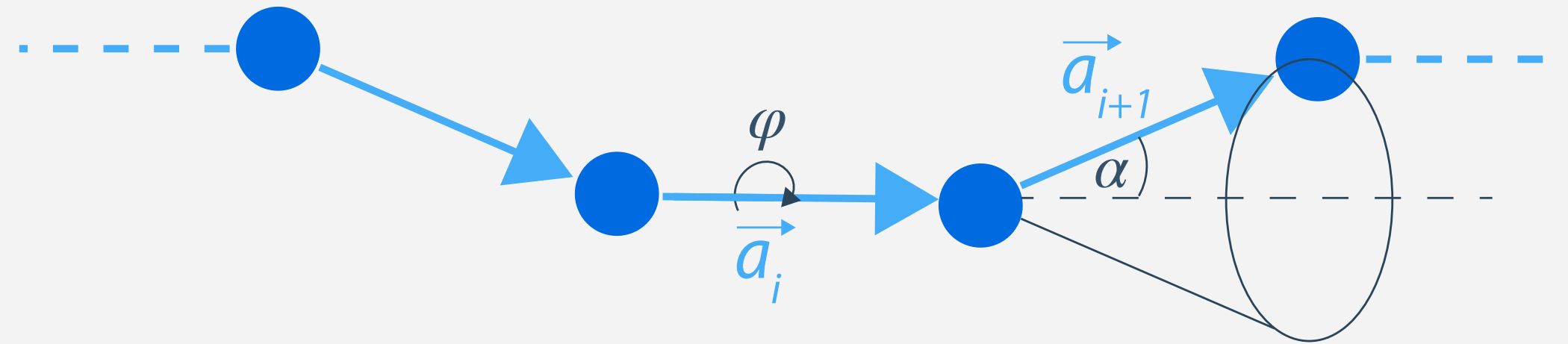
Where is the most probable position of any shot?

What would your average points be?



Freely Rotating Chain Model

- all torsion angles $-\pi < \varphi \leq \pi$ are assumed to be equally probable; tetrahedral angles α are fixed.



average projection from \vec{a}_{i+1} on \vec{a}_i : $|a| \cos \alpha$

average projection from \vec{a}_{j+1} on \vec{a}_i : $|a| \cos^{|j-i|} \alpha$

mean square end-to-end distance:

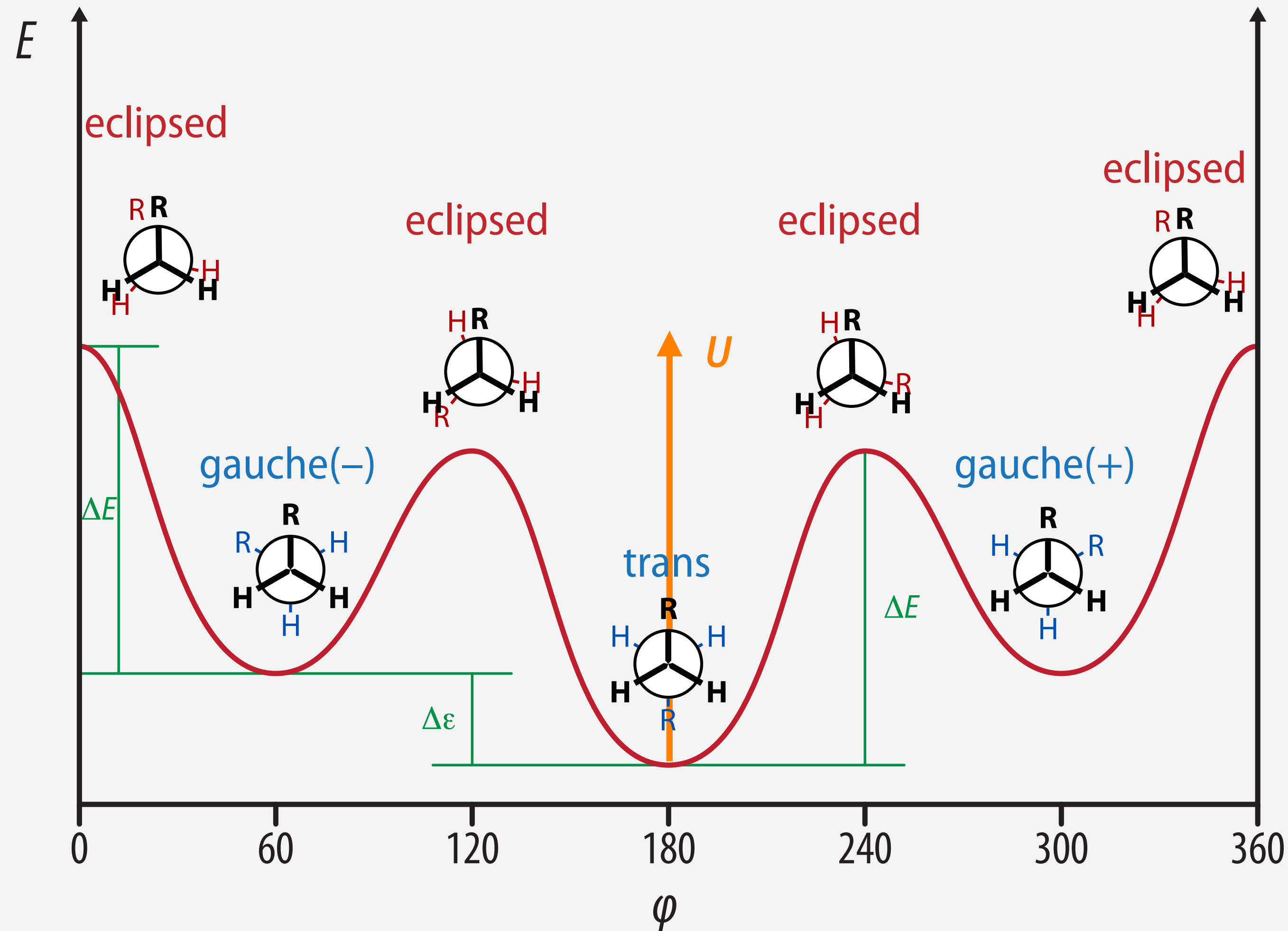
$$\langle R_n^2 \rangle = n |a|^2 + \langle \sum_{i \neq j} \vec{a}_i \cdot \vec{a}_j \rangle = n |a|^2 + 2 |a|^2 \sum_{i < j}^n \cos^{j-i} \alpha$$

- dependence of $\langle R^2 \rangle$ of an ideal linear chain on number of bonds, bond length, and bond angle:

$$\text{for large } n: \langle R_n^2 \rangle = na^2 \left(\frac{1 + \cos \alpha}{1 - \cos \alpha} \right)$$

Hindered Rotation Model

- constant bond lengths and angles, independent torsion angles with rotation hindered by a potential, U



for large n :

$$\langle R^2 \rangle = na^2 \left(\frac{1 + \cos\alpha}{1 - \cos\alpha} \right) \left(\frac{1 + \langle \cos\varphi \rangle}{1 - \langle \cos\varphi \rangle} \right)$$

Chain Flexibility and the Characteristic Ratio

- the mean-square end-to-end distance can be approximated for long chains:

Flory's characteristic ratio

$$\langle R^2 \rangle \cong C_\infty n a^2$$

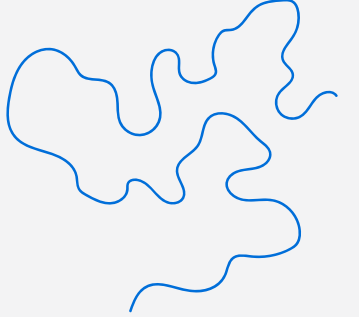
freely jointed chain: $C_\infty = 1$

freely rotating chain $C_\infty = \frac{1 + \cos\alpha}{1 - \cos\alpha}$

hindered rotation $C_\infty = \left(\frac{1 + \cos\alpha}{1 - \cos\alpha} \right) \left(\frac{1 + \langle \cos\varphi \rangle}{1 - \langle \cos\varphi \rangle} \right)$

- the characteristic ratio is a correction term for chain rigidity/flexibility.

random coil: 
 $C_\infty = 1$

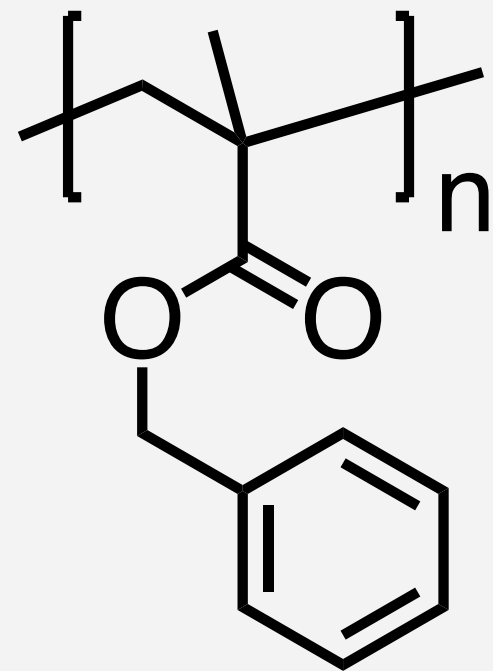
stretched conformation: 
 $C_\infty \gg 1$

$$\langle R^2 \rangle \cong \begin{cases} n \cdot a^2 & \text{freely jointed chain} \\ 2n \cdot a^2 & \text{freely rotating chain} \\ 3n \cdot a^2 & \text{hindered rotation} \end{cases}$$

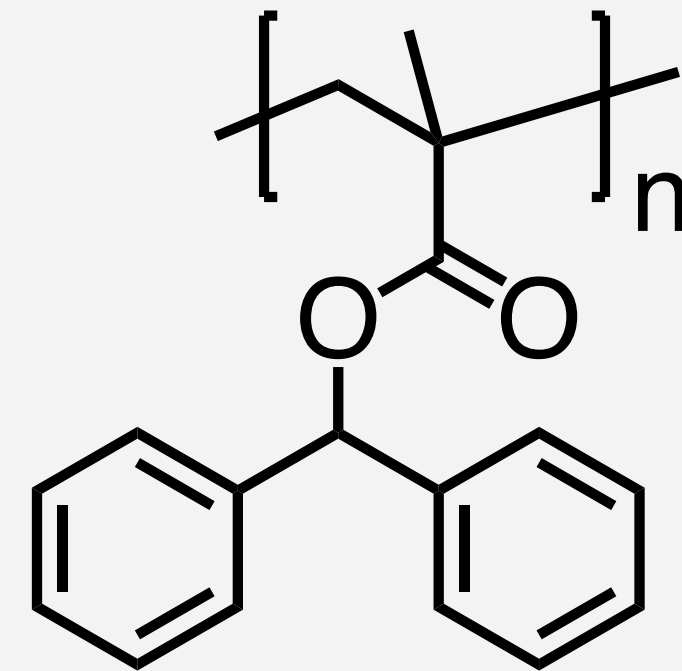
↓ chain rigidity

Examples

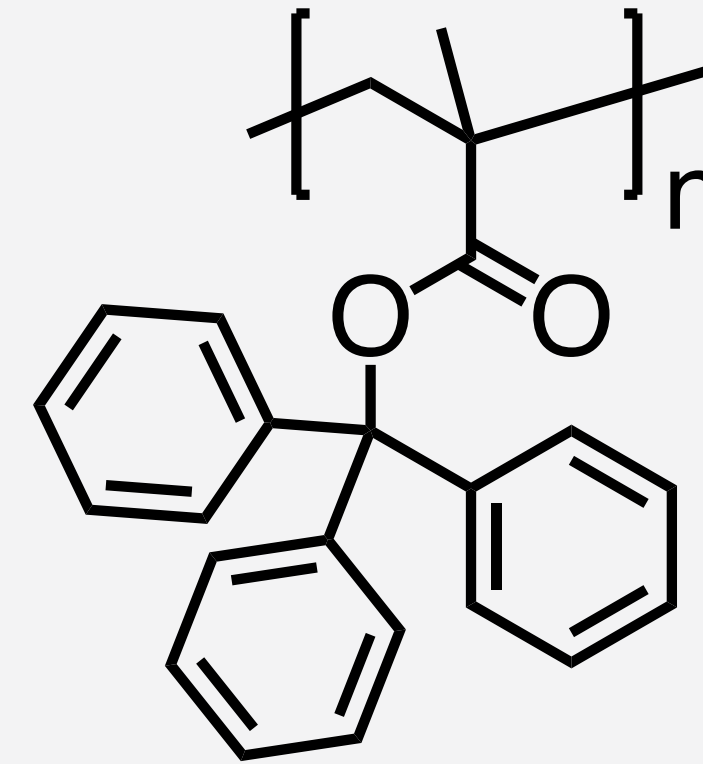
several
poly(methylmethacrylate)s



$$C_{\infty} = 10$$

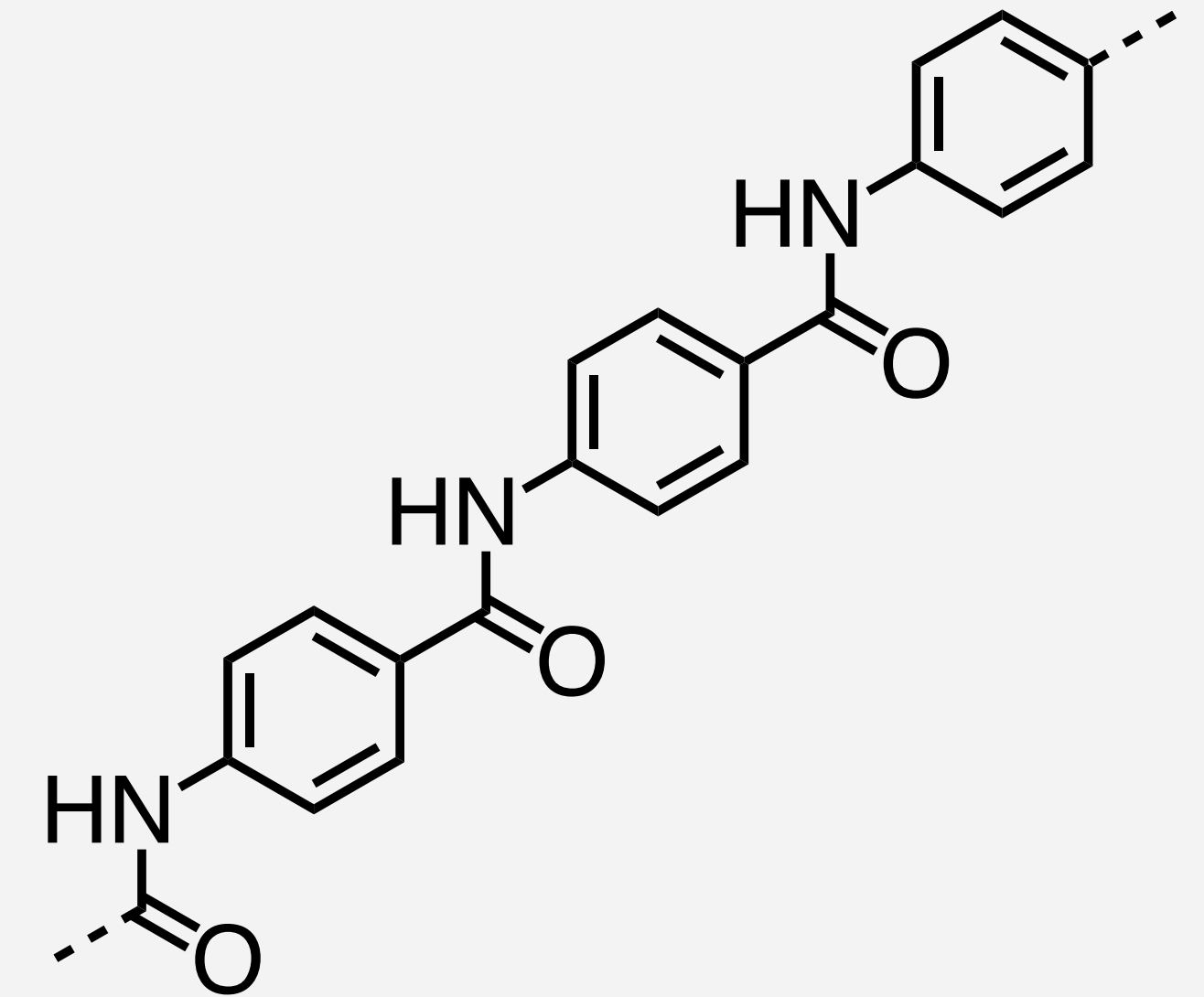


$$C_{\infty} = 14$$



$$C_{\infty} = 20.3$$

poly(*p*-benzamide)



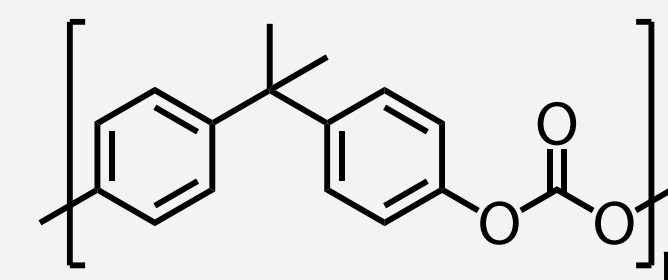
$$C_{\infty} = 325 \text{ why?}$$

- chain stiffness increases with increasing side chain bulkiness (limited rotation around main chain bonds)
- aromatic rings, double bonds increase the rigidity of a polymer chain

Typical Values of C_∞ in Solution

Polymer	Solvent	T [°C]	C_∞
polyethylene	1-dodecanol	138	6.7
polystyrene (atactic)	cyclohexane	35	10.2
polypropylene (atactic)	cyclohexane	92	6.8
polyisobutylene	benzene	24	6.6
polyvinylacetate	<i>i</i> -pentanone + hexane	25	8.9
polyoxomethylene	aqueous K_2SO_4	35	4.0
polycarbonate	methylenechloride	25	2.2
poly(benzobisoxazole)			93
poly(<i>p</i> -benzamide)			325

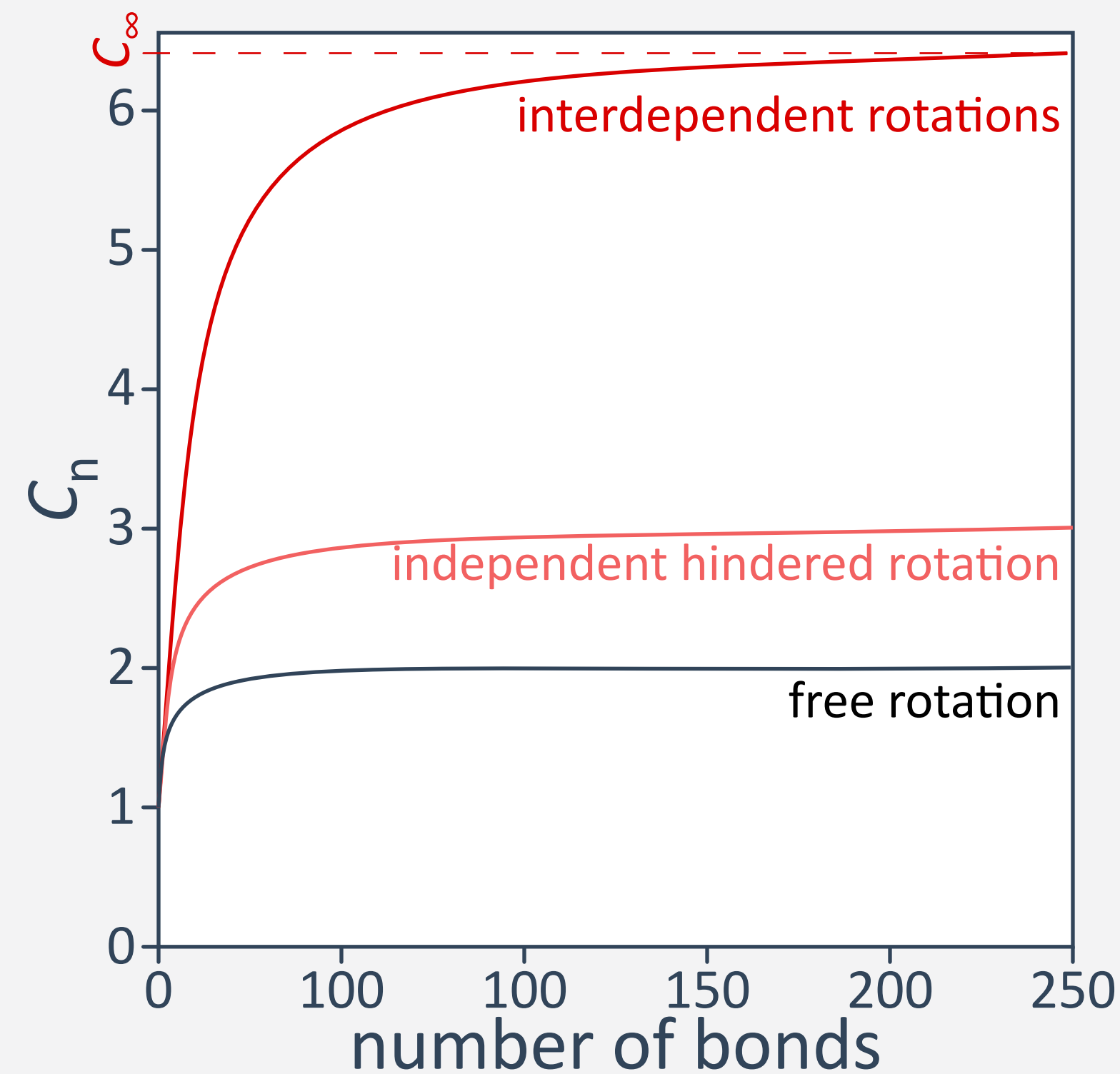
- be careful with an interpretation (see **Exercise Sheet**):



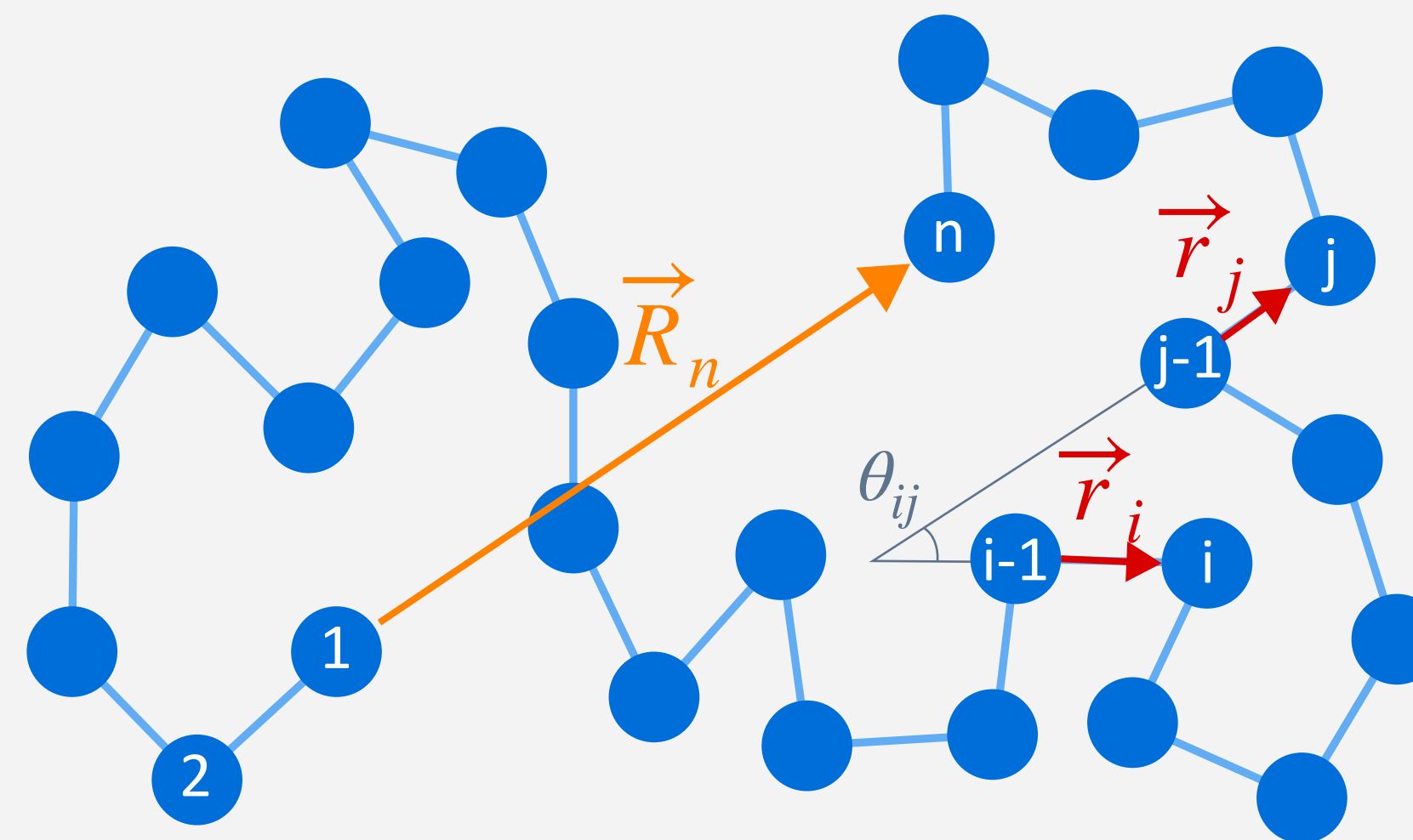
$$C_\infty = 2.2?$$

Deviations from Ideal-Chain Behaviour in PE

- Flory's characteristic ratio, C_n , approaches a finite value only for $n \rightarrow \infty$ ($C_\infty = 6.7$ for PE)



$$C_n = \frac{1}{n} \sum_{i=1}^n C'_i \quad C'_i \equiv \sum_{j=1}^n \langle \cos \theta_{ij} \rangle$$



local and non-local, more far reaching interactions

- real chains: further correlations between bond vectors due to forces acting on individual chain elements

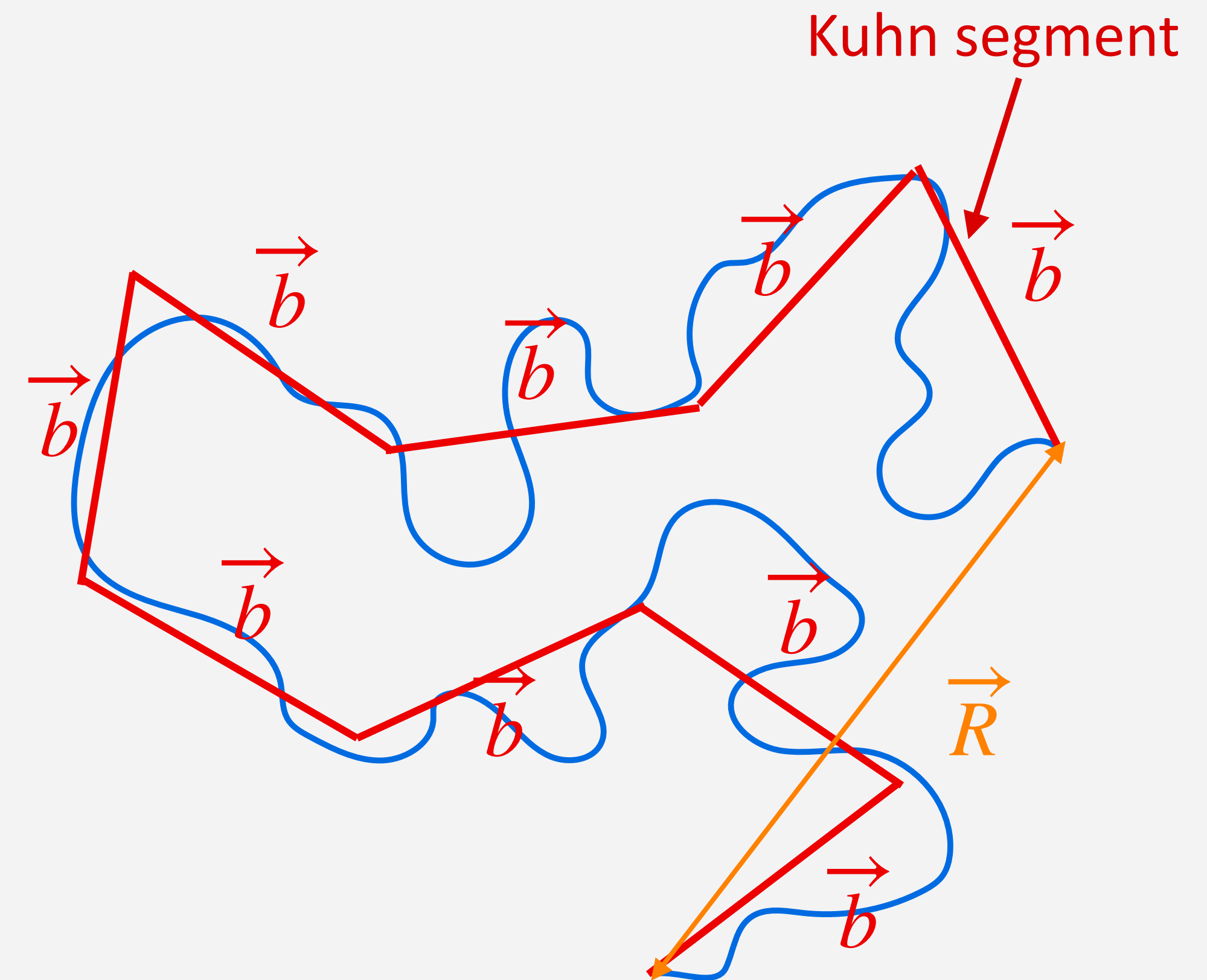
Kuhn Segments

- real polymer chains can be represented by an equivalent freely-jointed chain:

same projection length:
(maximum end-to-end distance) $na = Nb$

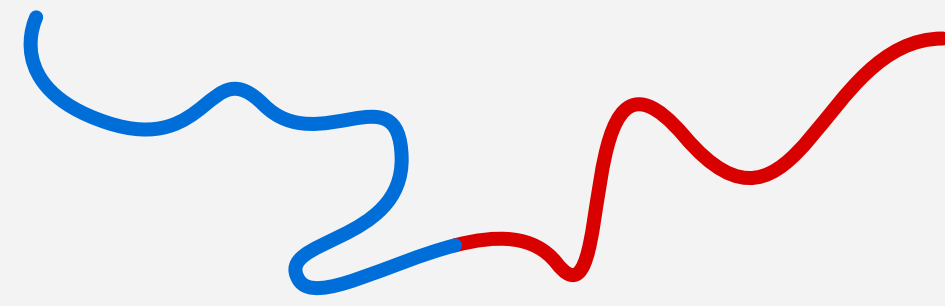
same end-to-end distance: $\langle R^2 \rangle = Nb^2 = C_\infty na^2$

$$b = \frac{C_\infty na^2}{R_{max}} = aC_\infty$$



- b is the **Kuhn segment length** which increases with chain stiffness (see also **Exercise**)

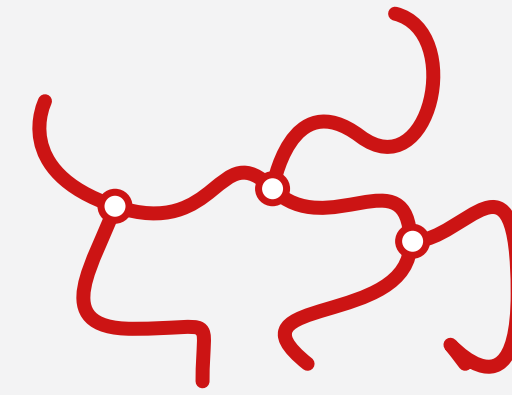
End-to-End Distance: Not Always An Appropriate Measure



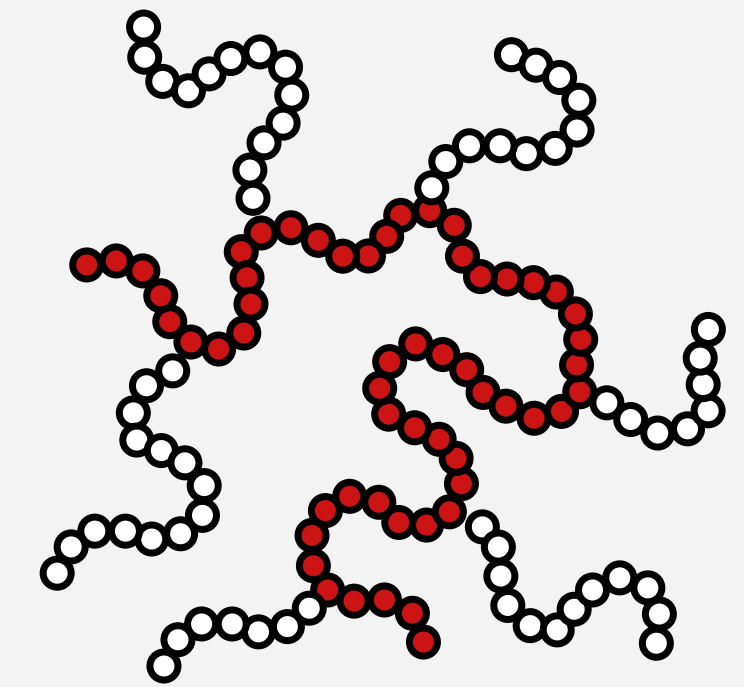
block copolymer



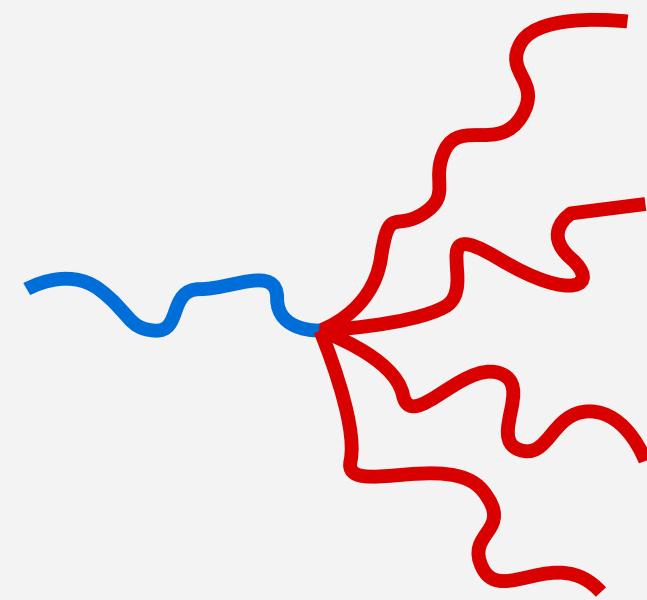
star polymer



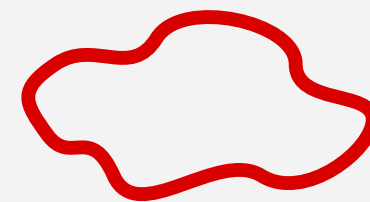
branched polymer



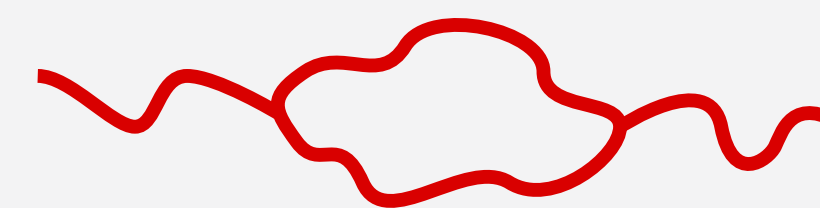
grafted polymer



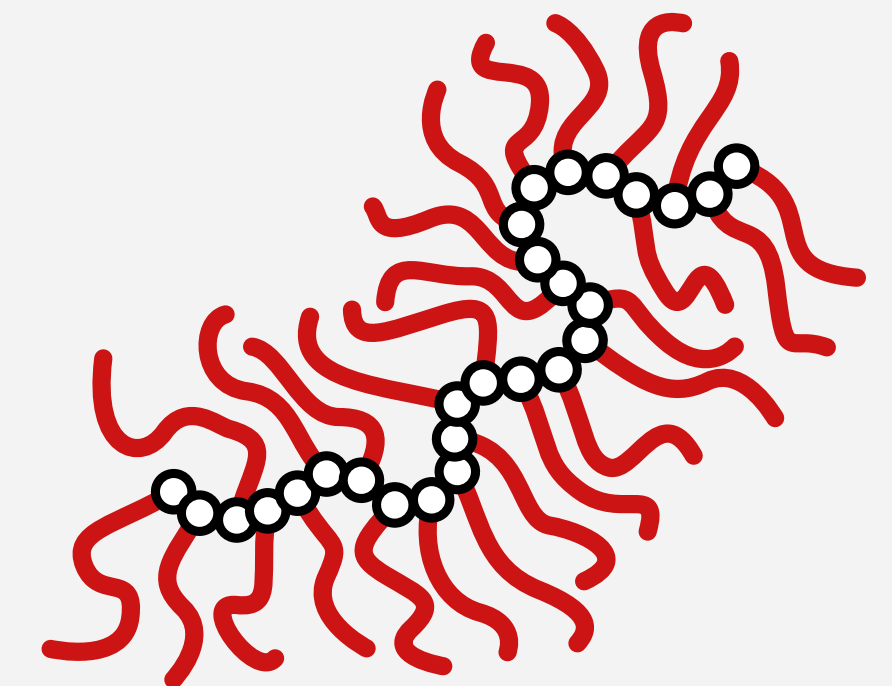
palm-tree AB_n



ring



coil-cycle-coil

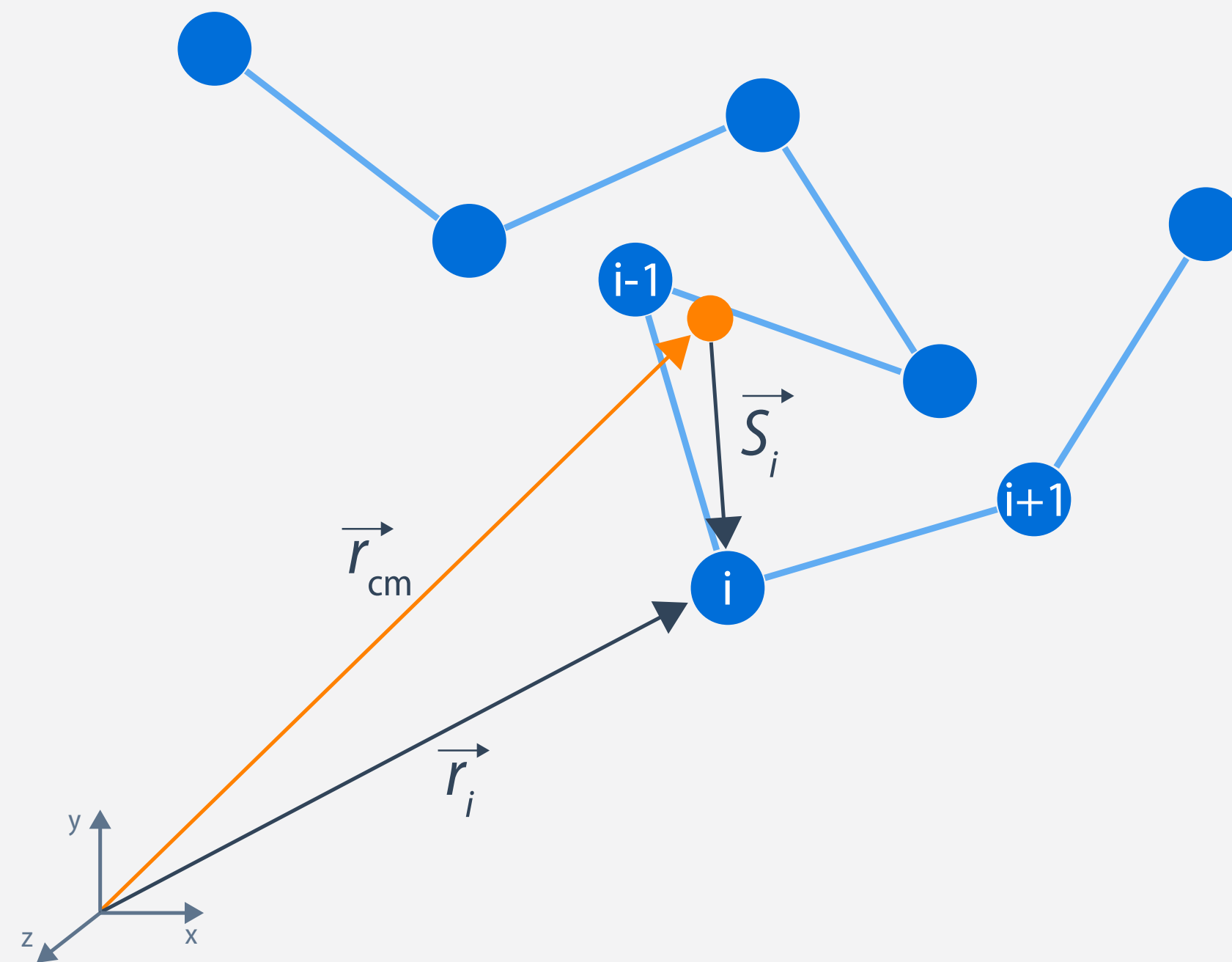


brush polymer

- the end-to-end distance of some polymers can not be defined unambiguously

Radius of Gyration

- **radius of gyration, R_g** , characterises the polymer size of any architecture (including branched or rings)



$$R_g^2 = \frac{1}{N} \sum_{i=1}^N (\vec{r}_i - \vec{r}_{\text{cm}})^2 = \frac{1}{N} \sum_{i=1}^N S_i^2$$

$$\sum_1^N \vec{S}_i = 0$$

- the mean-square of R_g relates to the mean-squared end-to-end distance for an ideal linear chain:

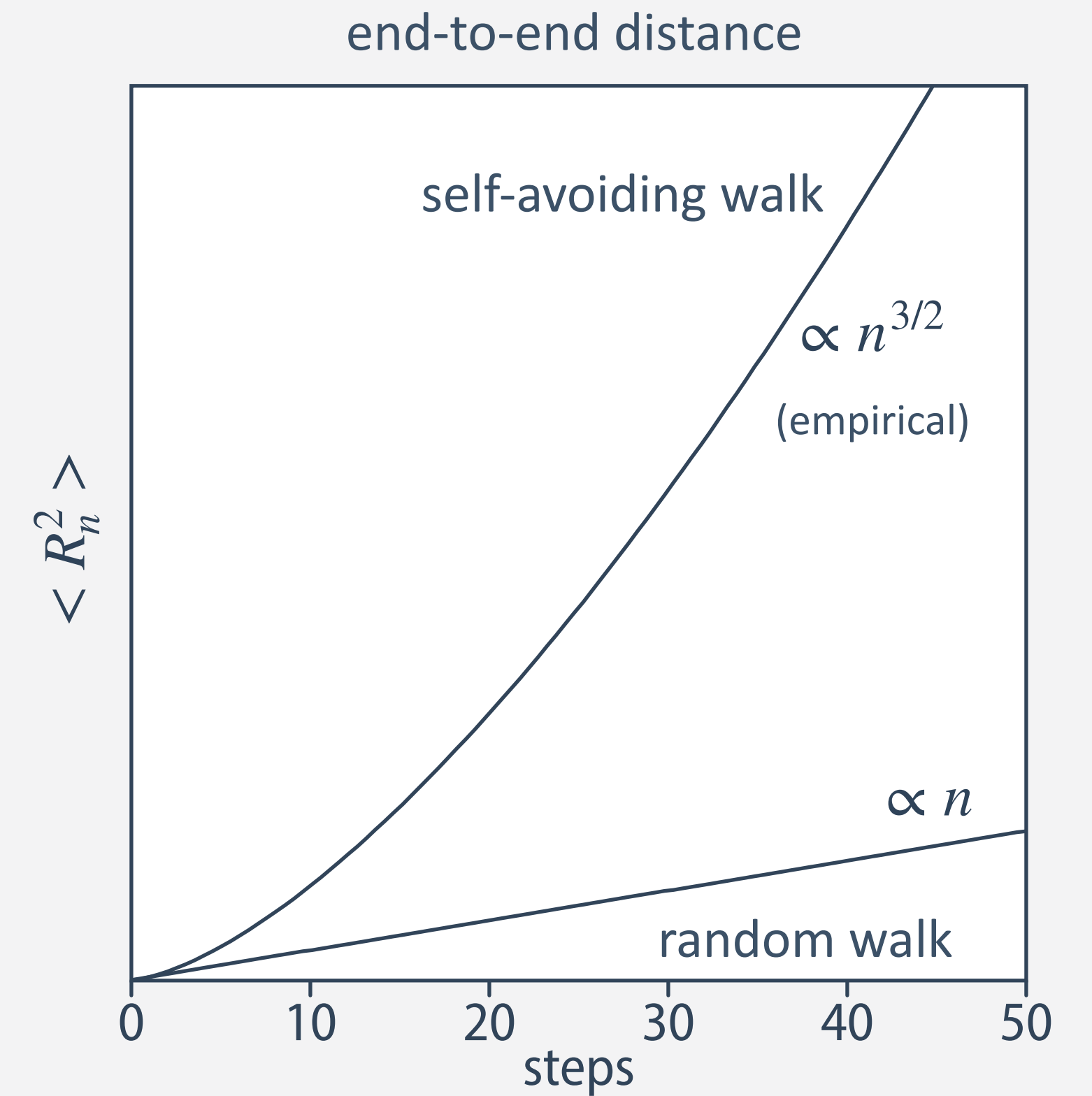
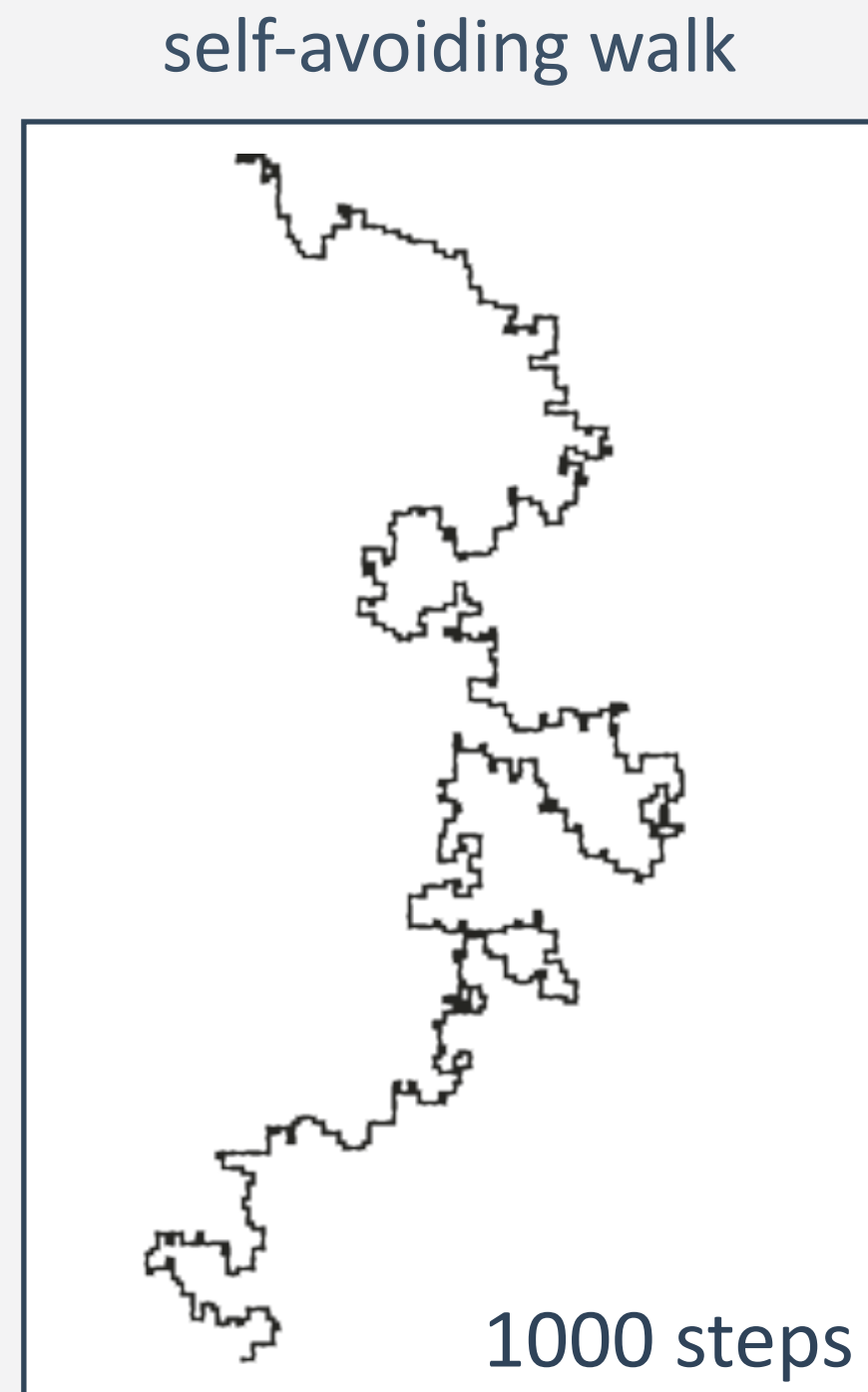
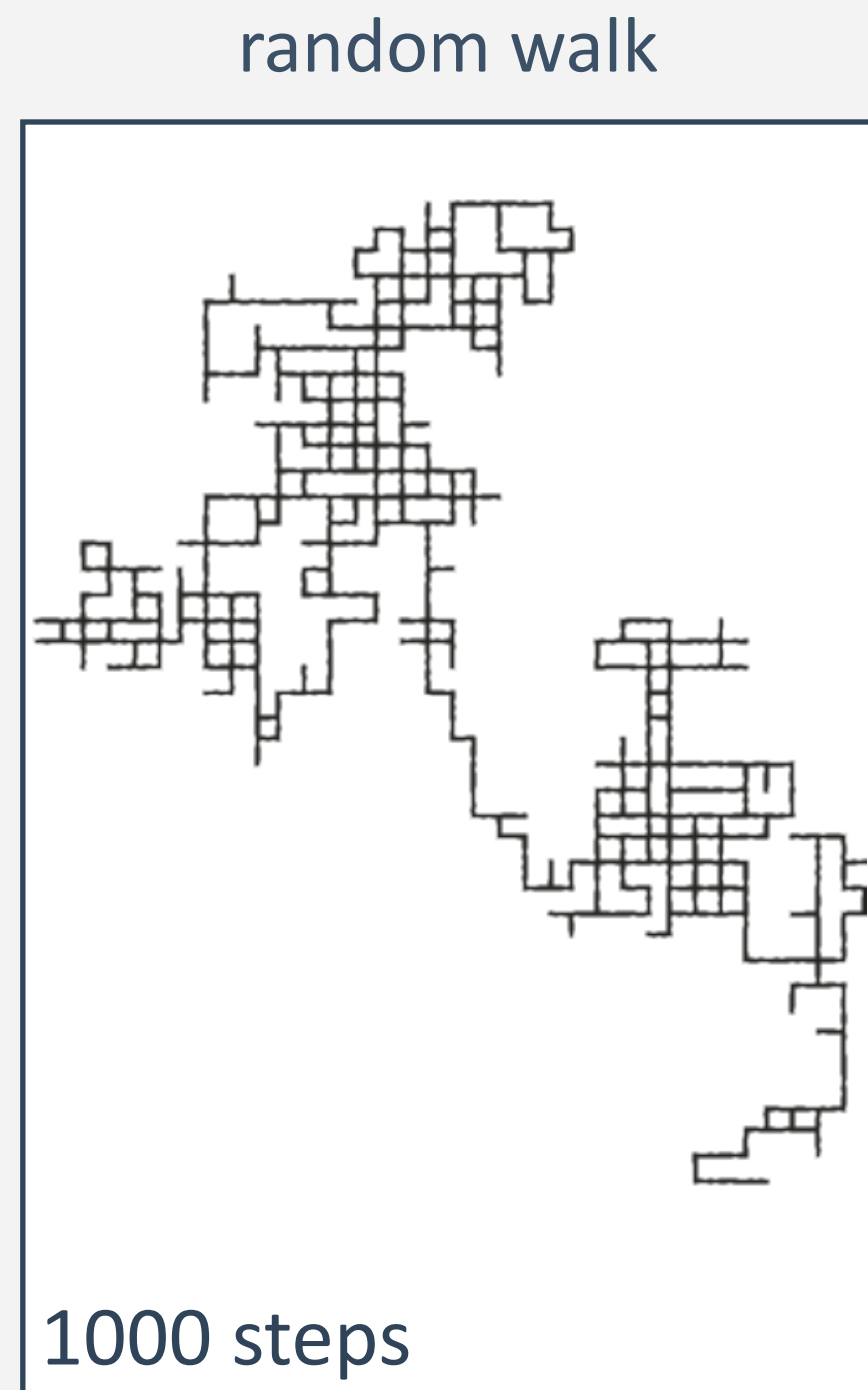
$$\text{for large } n: \langle R_g^2 \rangle = \frac{nl^2}{6} = \frac{\langle R_n^2 \rangle}{6}$$

2.2

Real Polymer Chains

Self-Avoiding Random Walk

- real chain: chain segments have a finite volume and they undergo interactions with their surrounding

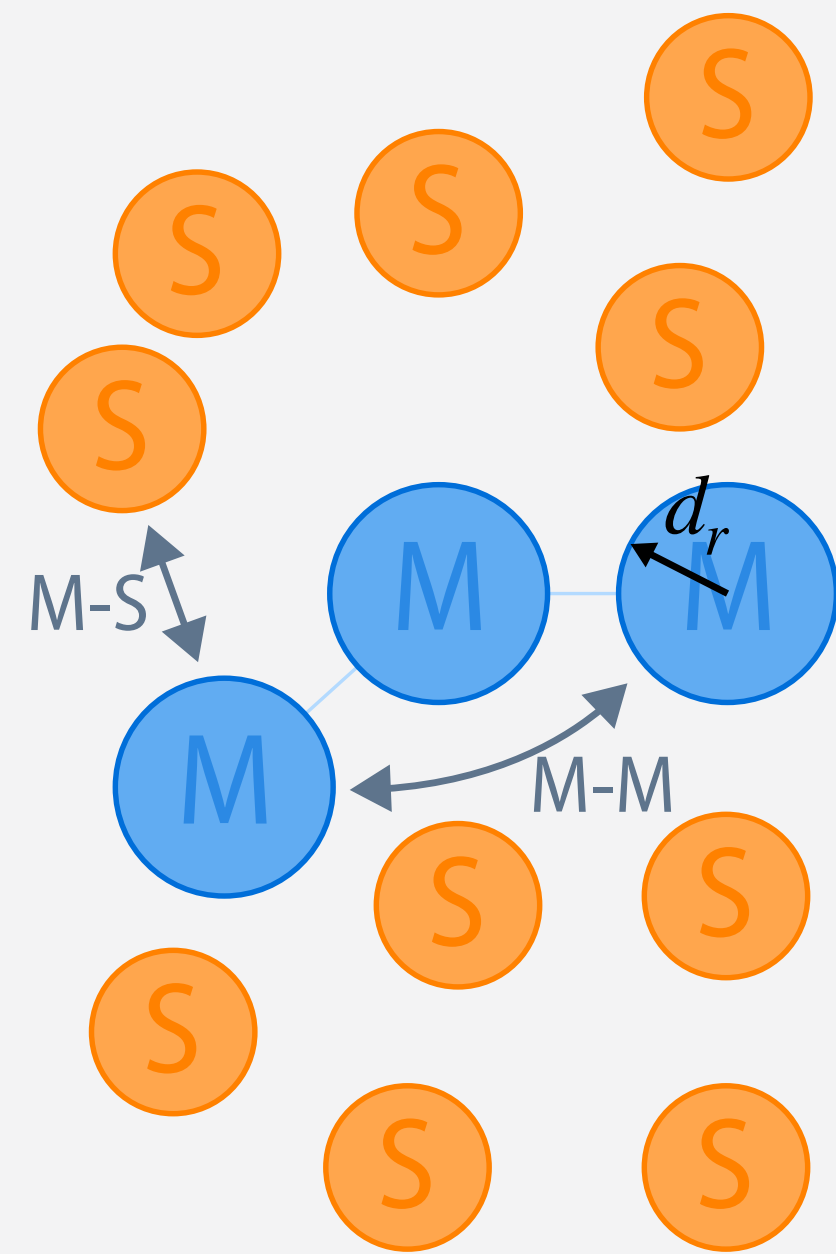


- real polymer chains may be mapped onto self-avoiding walks (volume excluded by other monomers)

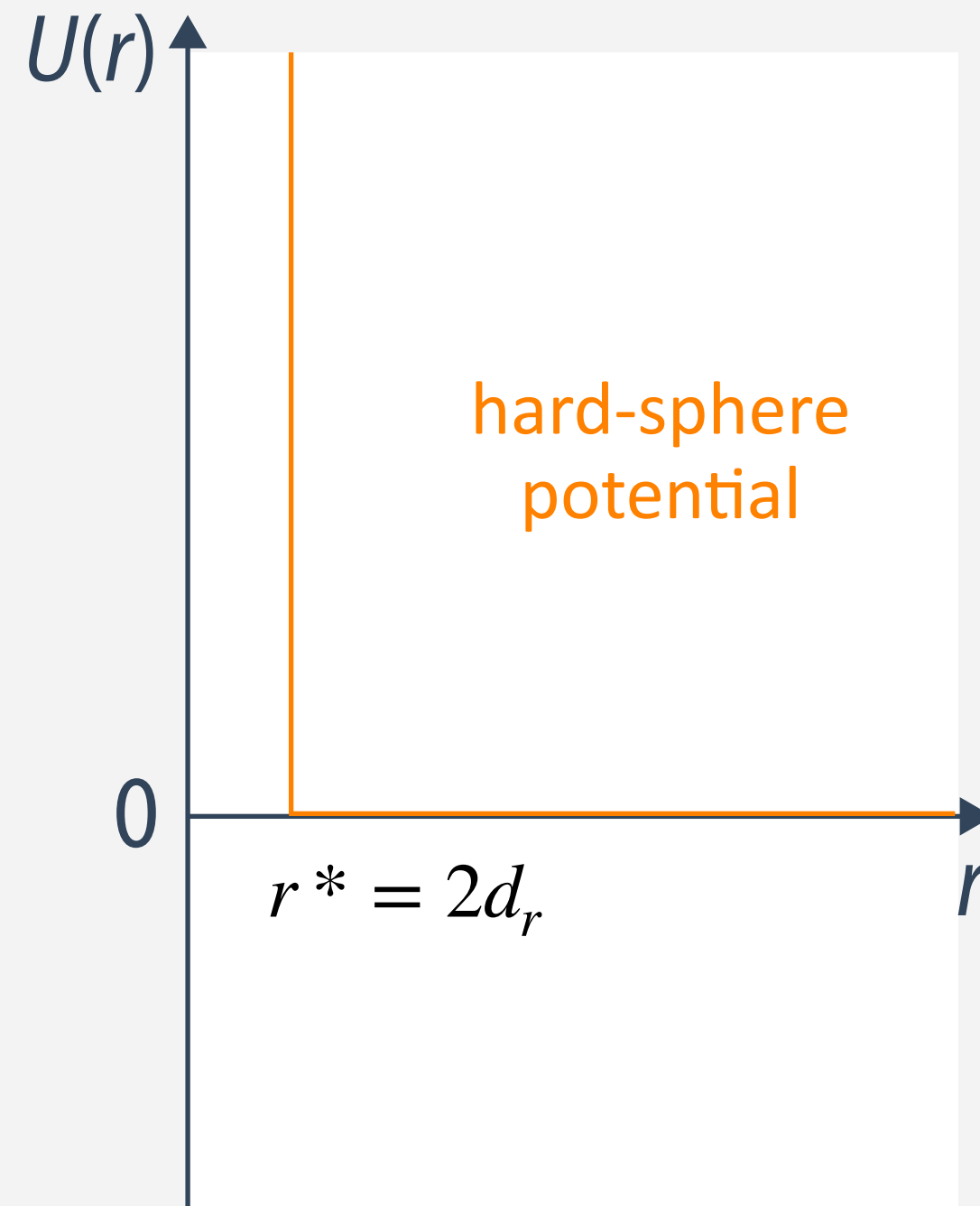
Effective Interaction Potentials

- polymer conformation is determined by monomer-monomer and monomer-solvent interactions

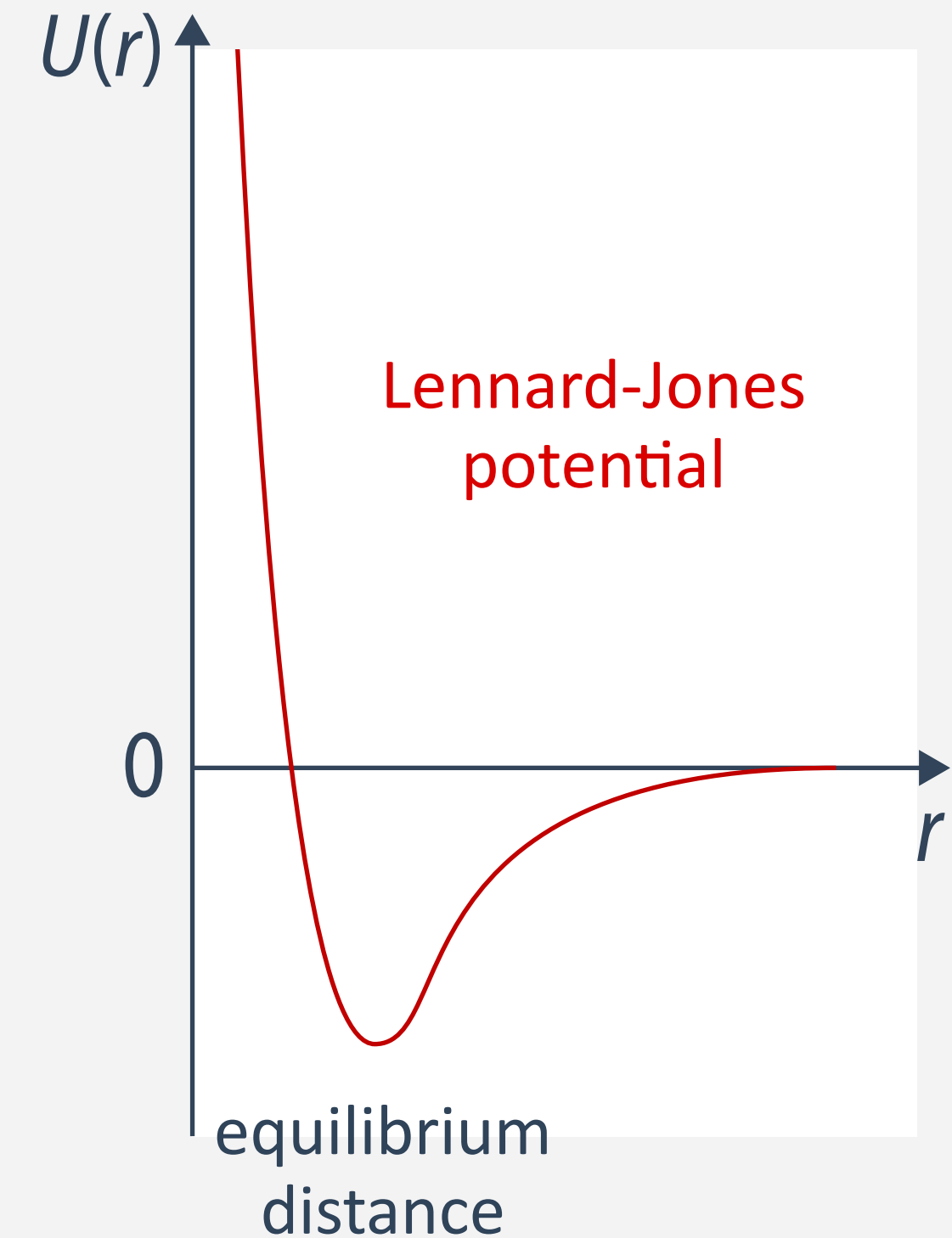
polymers in solution



zero monomer-monomer net interaction



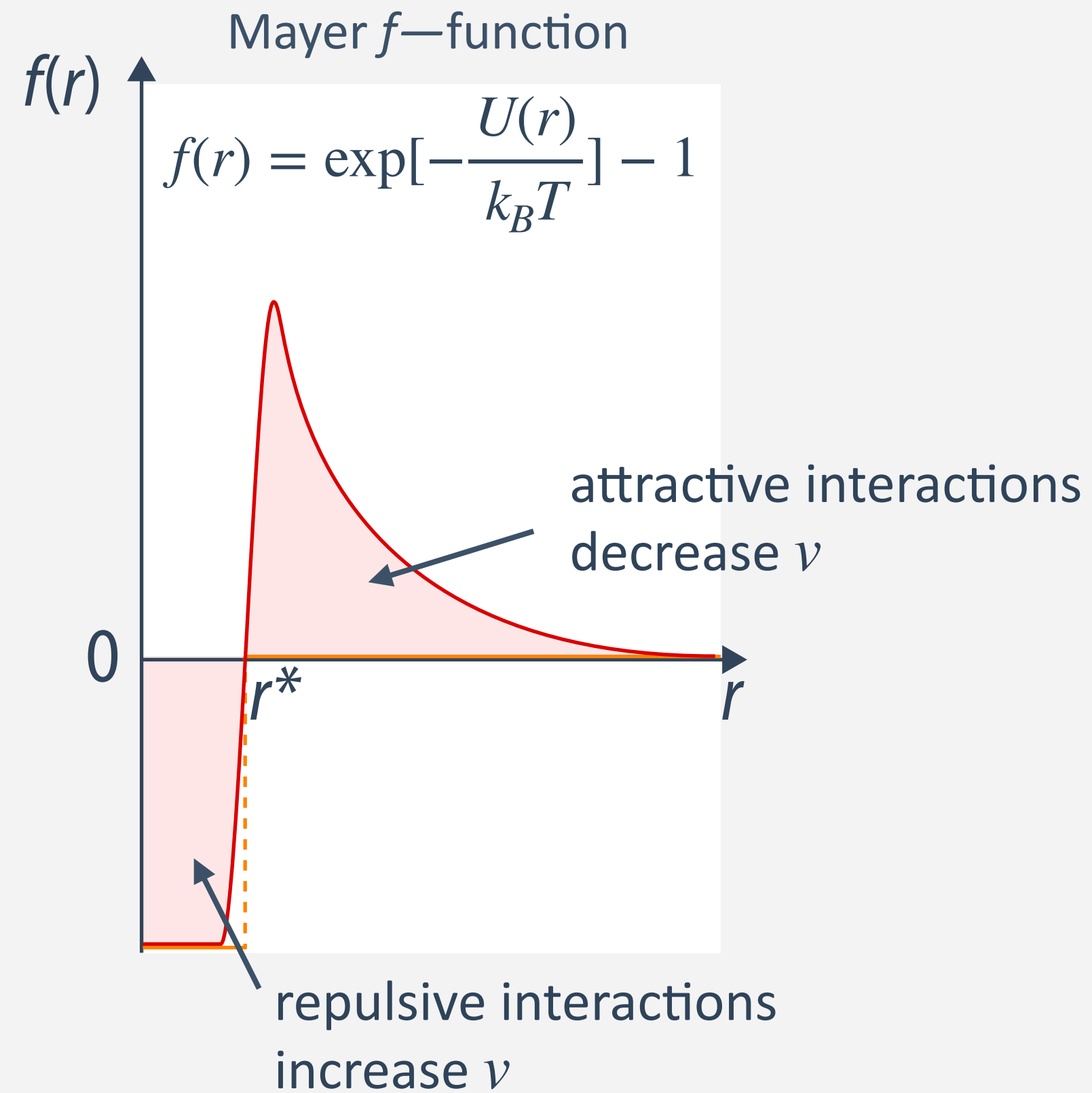
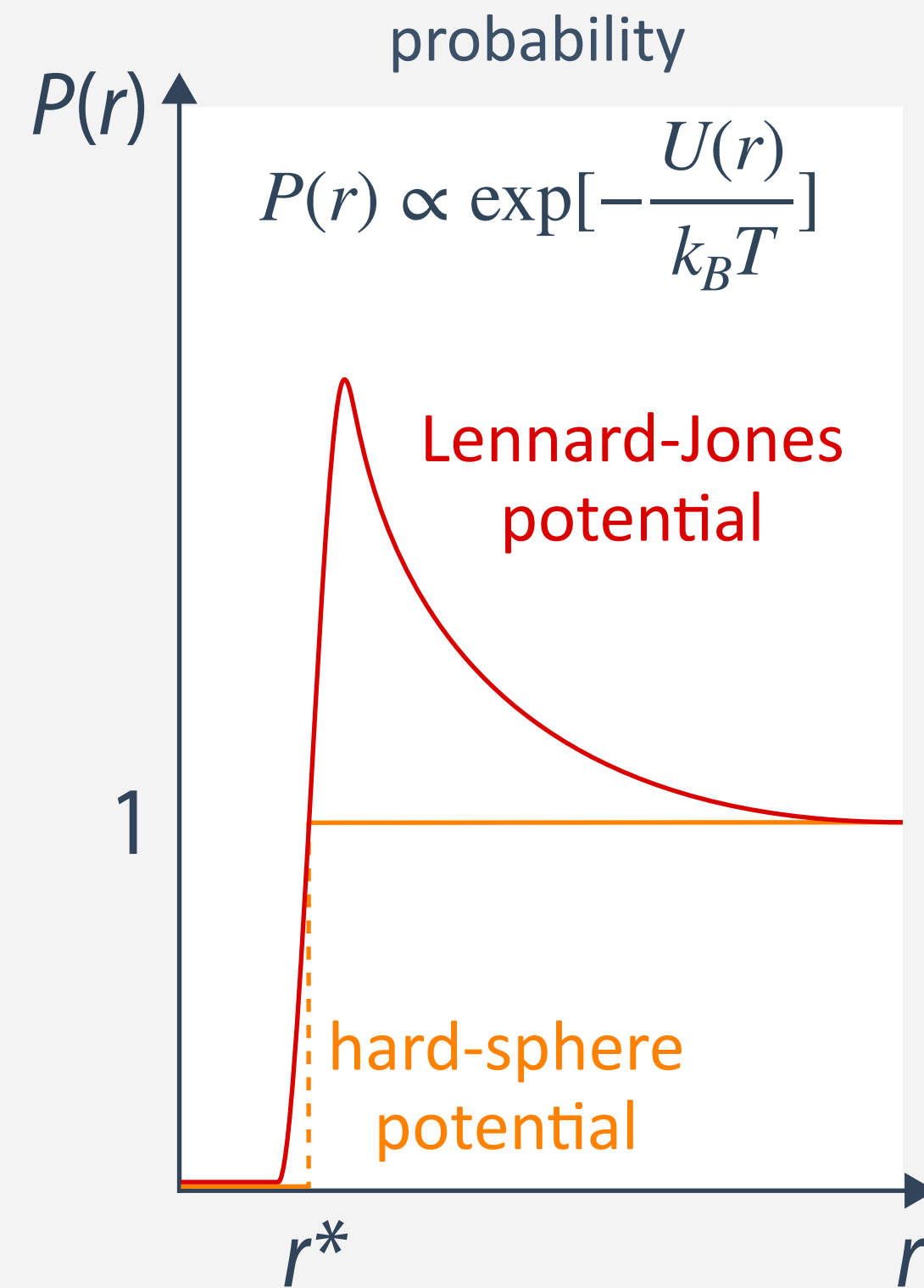
attractive monomer-monomer interaction



- Lennard Jones potential: monomer-monomer attraction, but strong repulsion at short distances

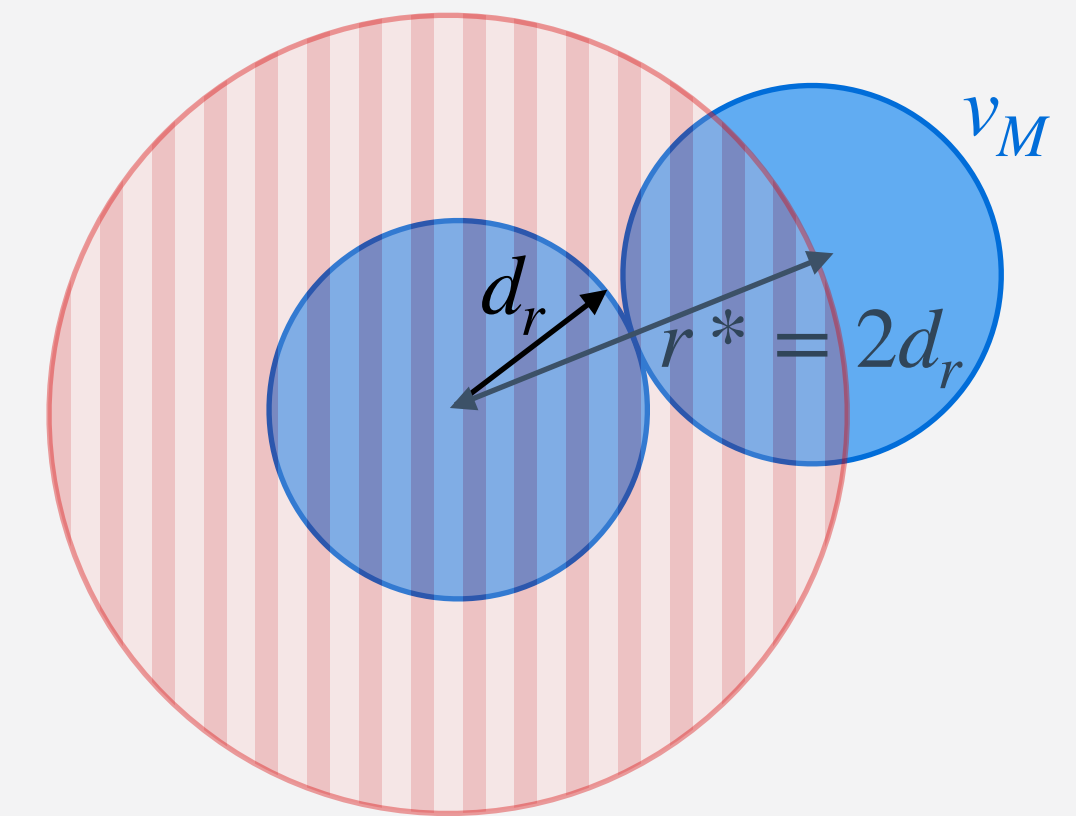
Excluded Volume

- probability of distance r between 2 monomers expressed by Boltzmann's distribution or Mayer f -function



excluded volume

$$\nu = - \int f(r) d^3 r$$



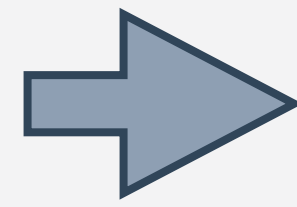
$$\nu_{max} = \frac{4\pi}{3} (2d_r)^3 = 8\nu_M$$

- excluded volume: space that each chain segment blocks to its surrounding

Classification of Solvents

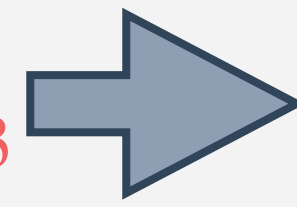
athermal

$$v = (2d_r)^3$$



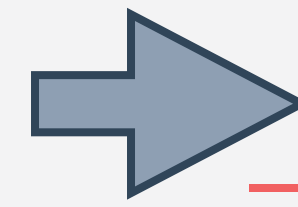
good

$$0 < v < (2d_r)^3$$



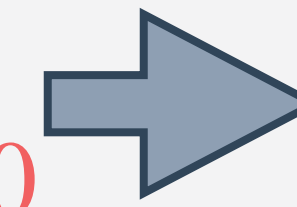
theta

$$v = 0$$



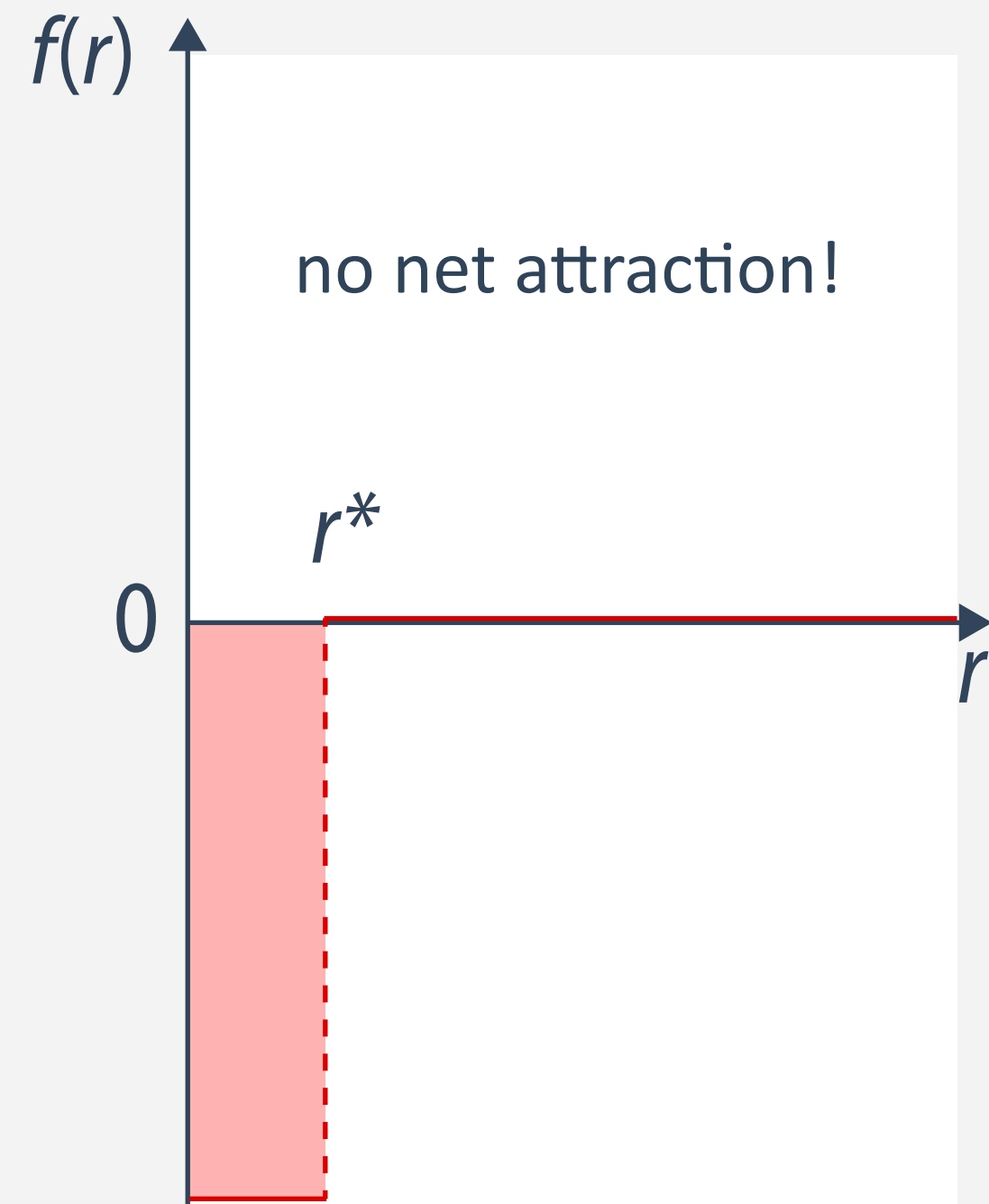
poor

$$-(2d_r)^3 < v < 0$$



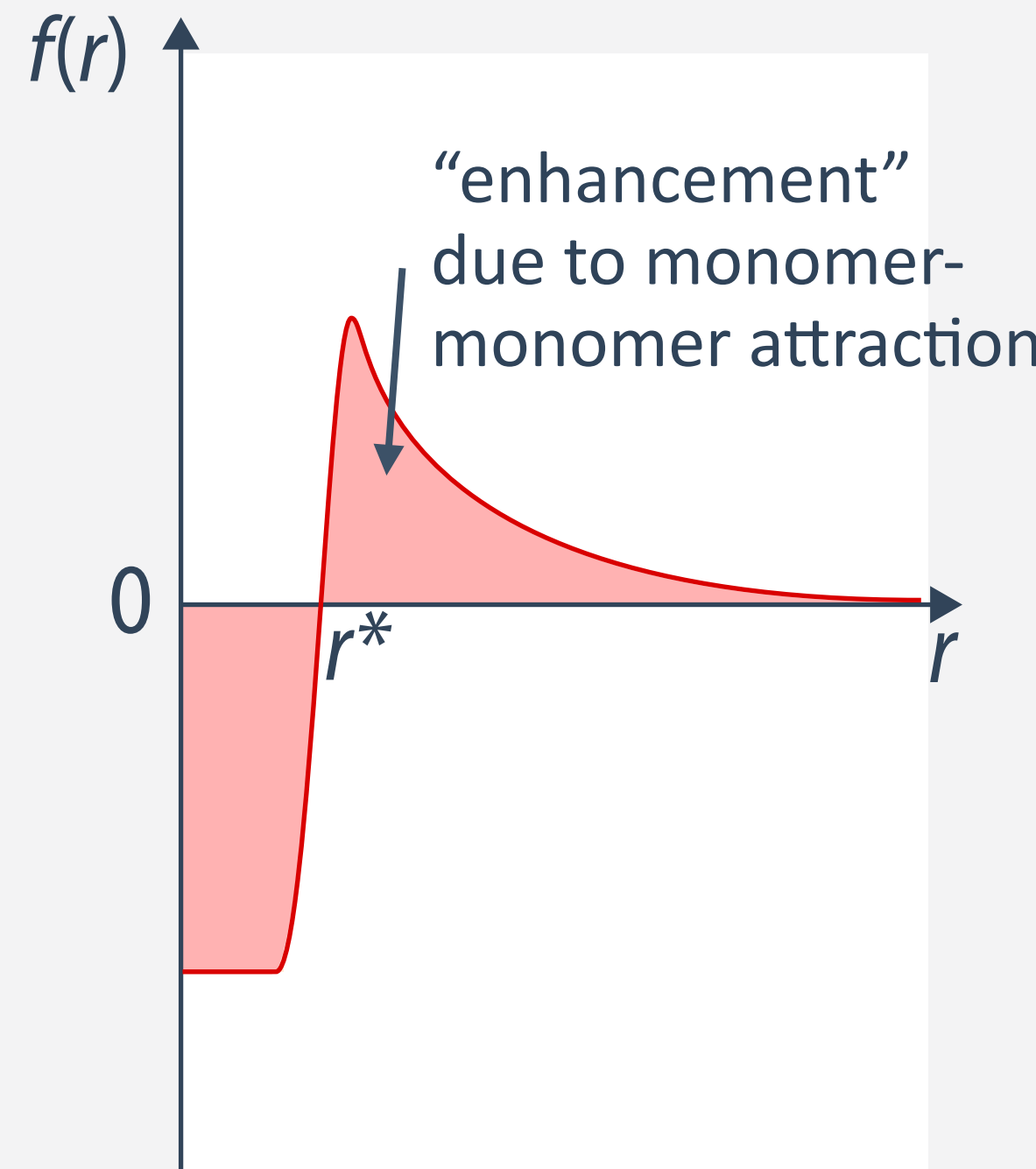
non-solvent

$$v = -(2d_r)^3$$



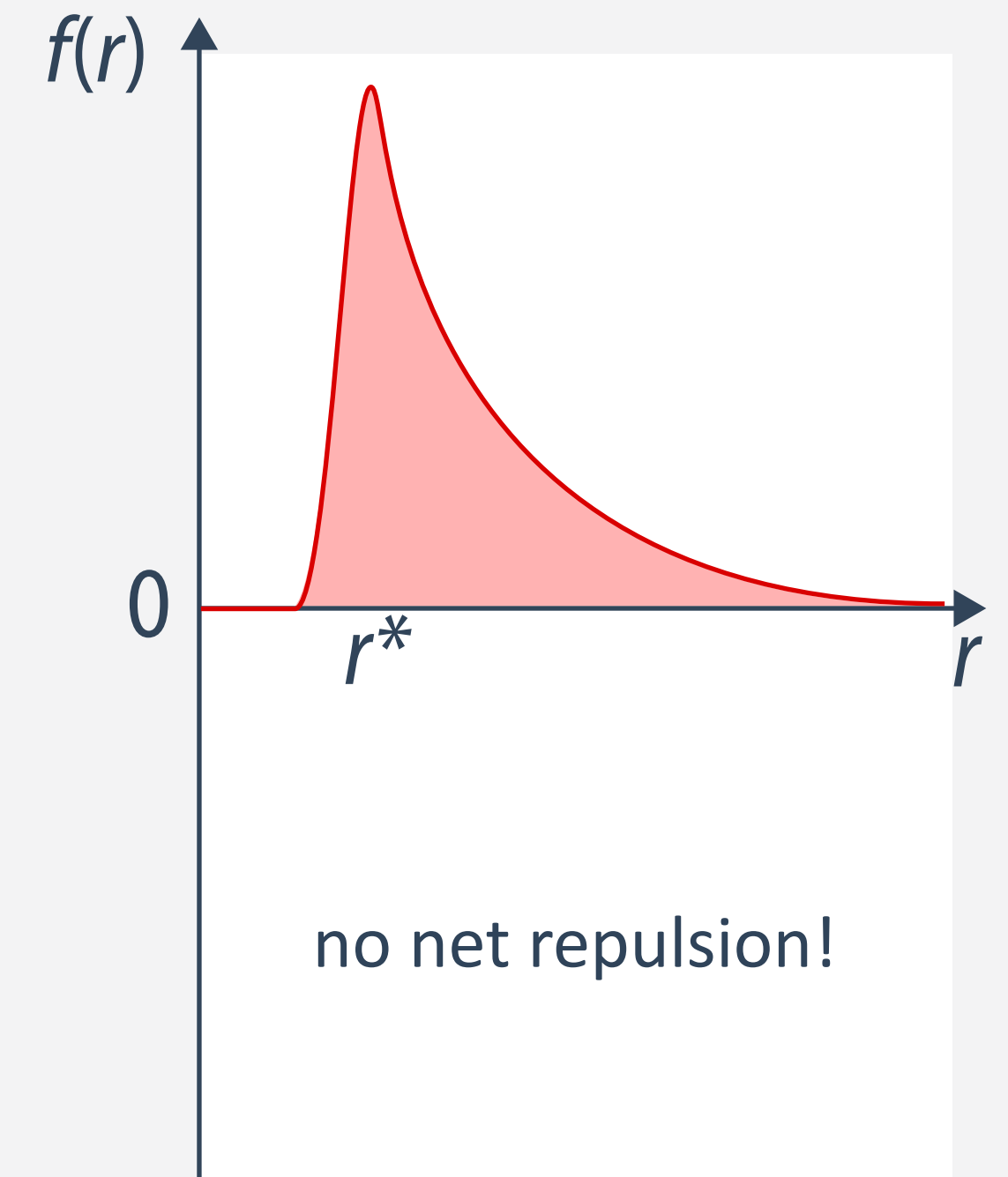
chains are swollen

ethyl benzene & PS



ideal chain behaviour

cyclohexane & PS ($T_\theta = 34.5 \text{ }^\circ\text{C}$)



chains are collapsed

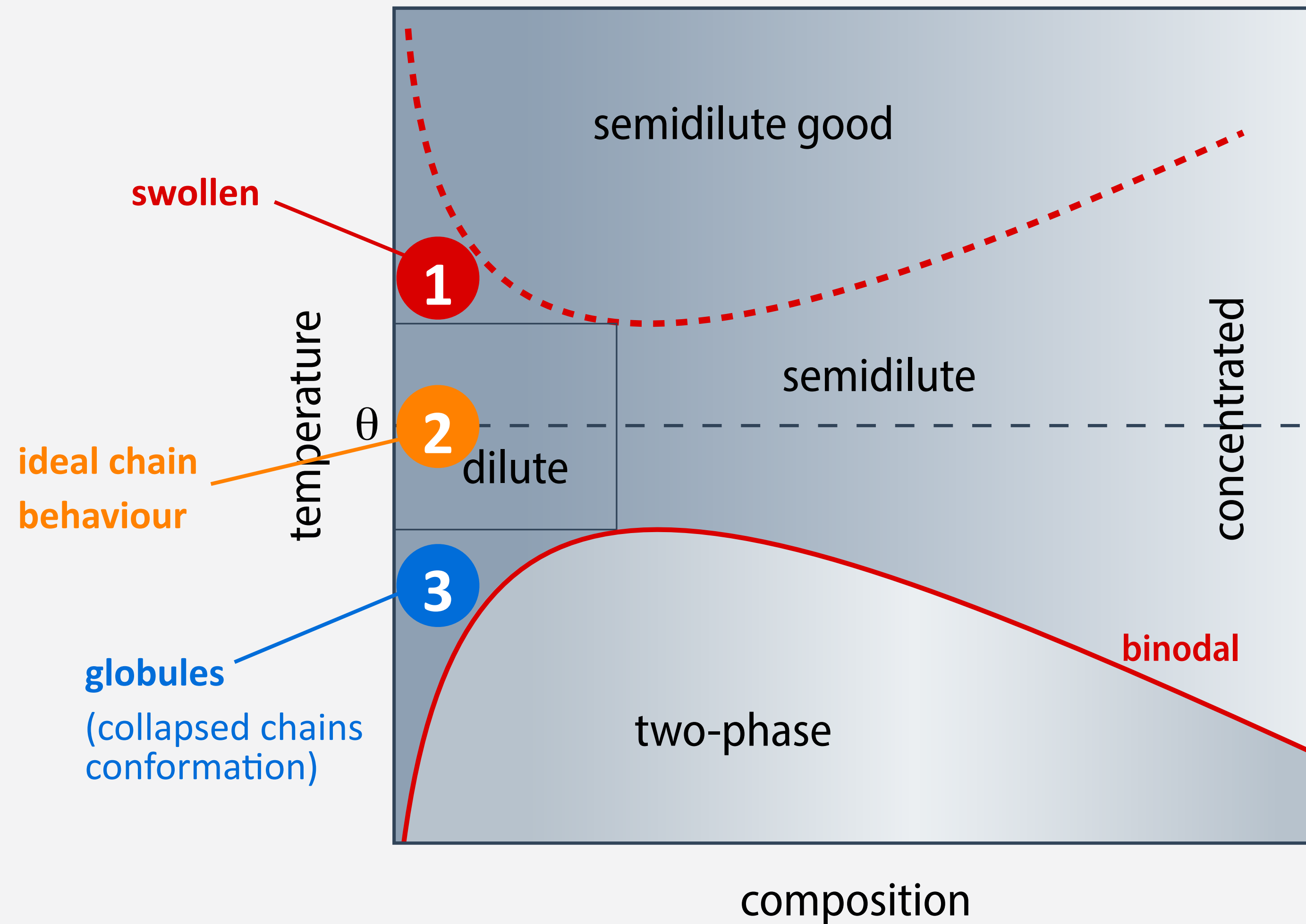
EtOH & PS

water & PS

- better solvent quality leads to polymer coil expansion, and a lower segmental density in the coil interior

Phase Diagram of Polymer Solutions

- at sufficiently low concentration, chains are well dispersed and do not phase-separate from the solvent

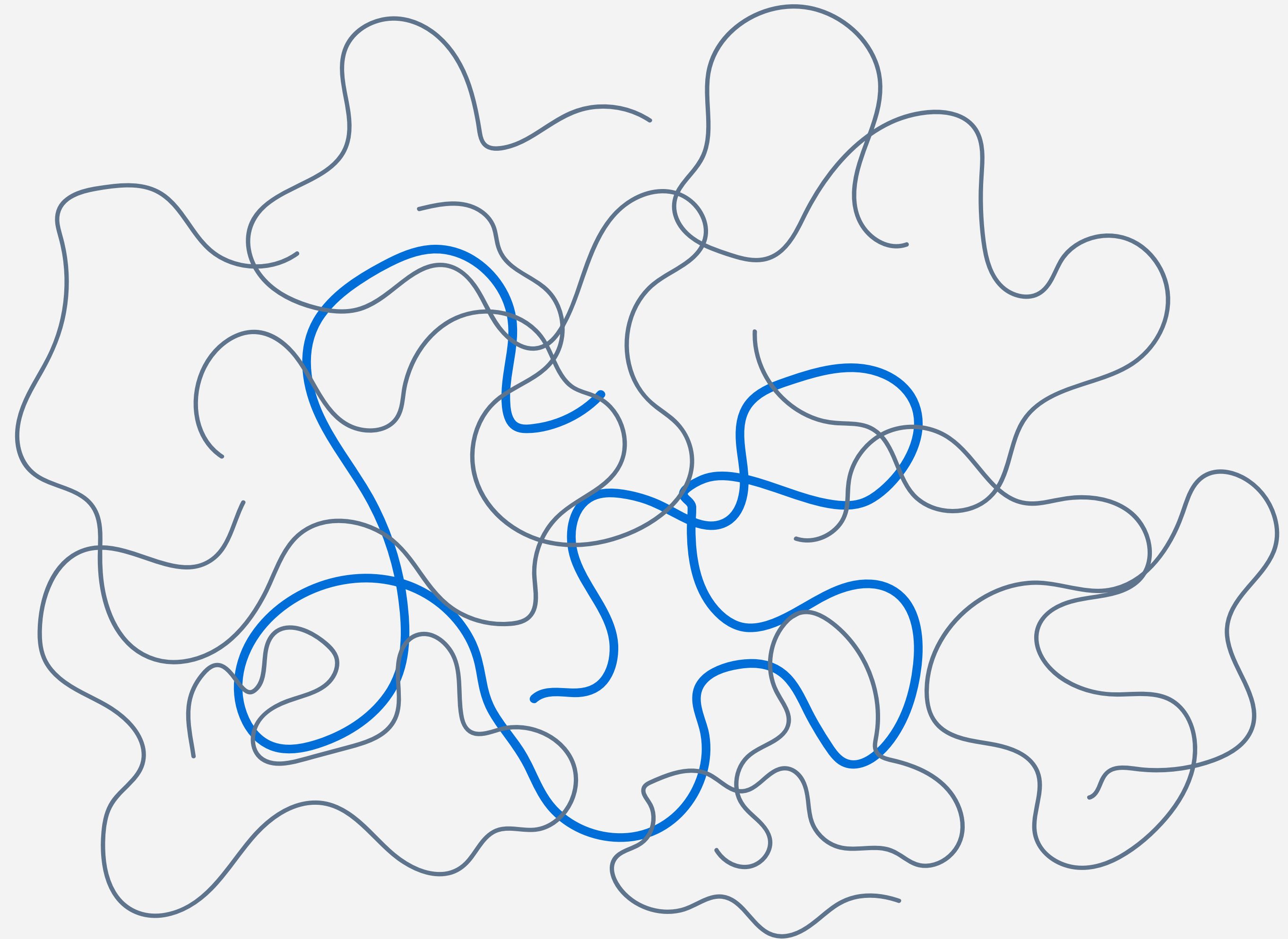


example:
cyclohexane & PS
 $T_{\theta} = 34.5 \text{ }^{\circ}\text{C}$

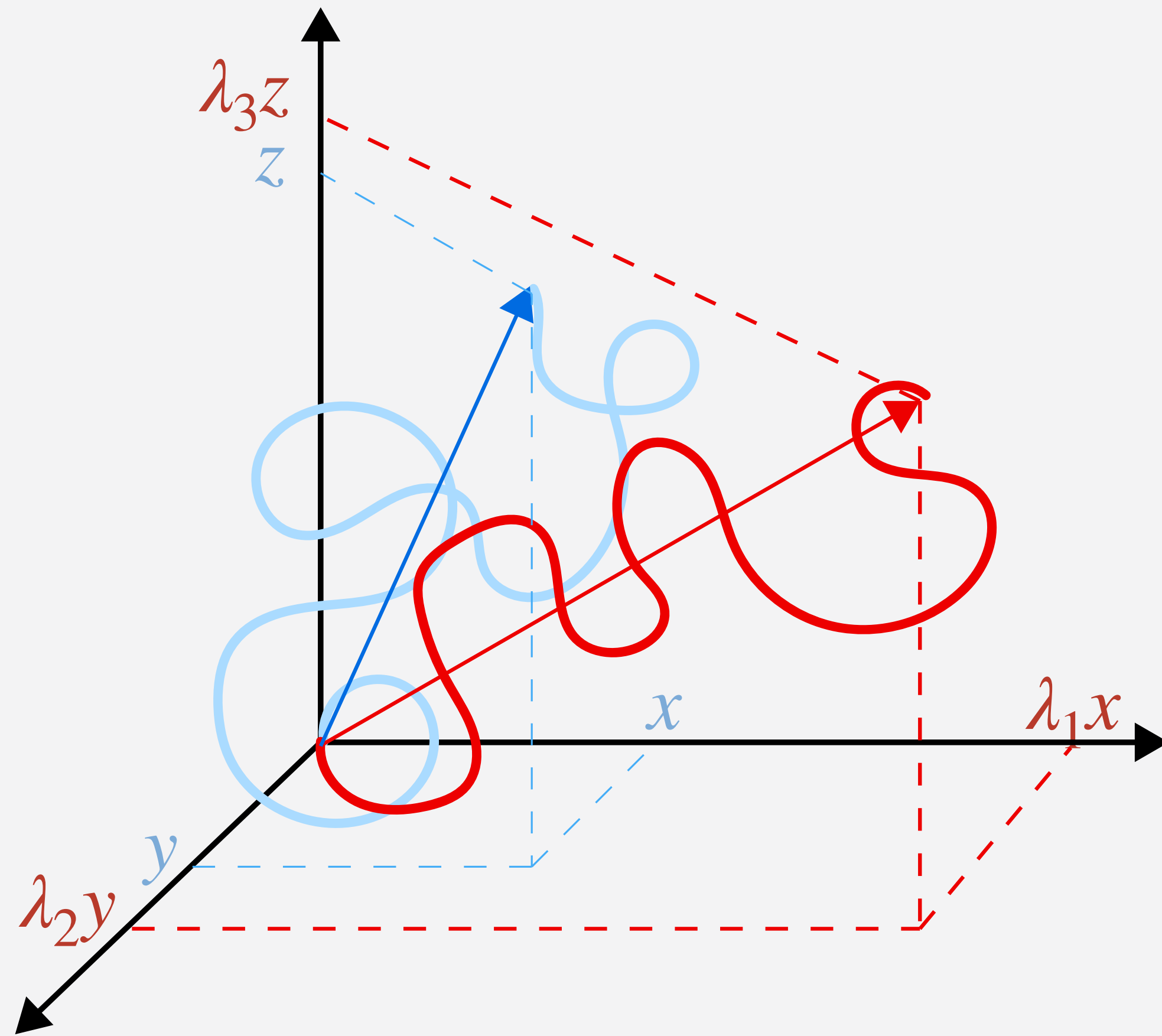
- at the θ -temperature, polymer behave ideal and are miscible with the solvent at any concentration

Omnipresence of the θ -State in Polymer Melts

- polymer melts constitute an athermal state (identical “monomer-monomer” & “monomer-solvent” interactions)
- same tendency of all chains to expand
- as a result, no chain can expand (excluded volume interactions are screened by overlapping segments of other coils)



- **polymer chains adopt their random coil conformation in polymer melts at any temperature!**



How does the end-to-end distance change upon deformation?

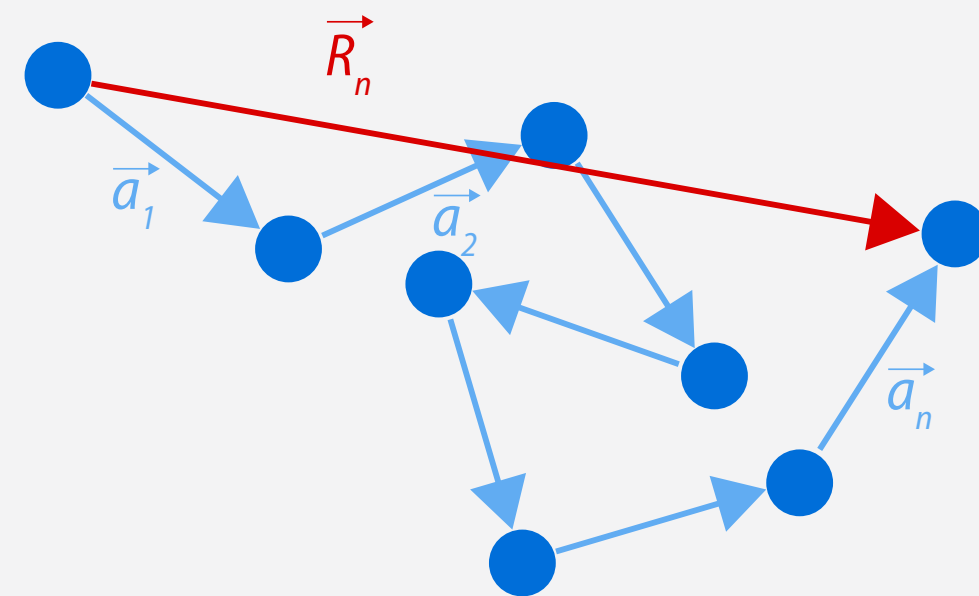
see Chapter 4.1 (Rubber Elasticity)

Learning Outcome

- every possible conformation of an ideal chain can be mapped onto a random walk
- a common feature to all ideal chain models is that size scales with \sqrt{M} for large n
- restrictions in available chain conformation with respect to the freely jointed chain is expressed using C_∞ , a measure of chain “stiffness”.
- accurate description of all polymer melts and certain behaviour in solution with the ideal chain model

$$\langle R_n^2 \rangle \cong C_\infty n a^2$$

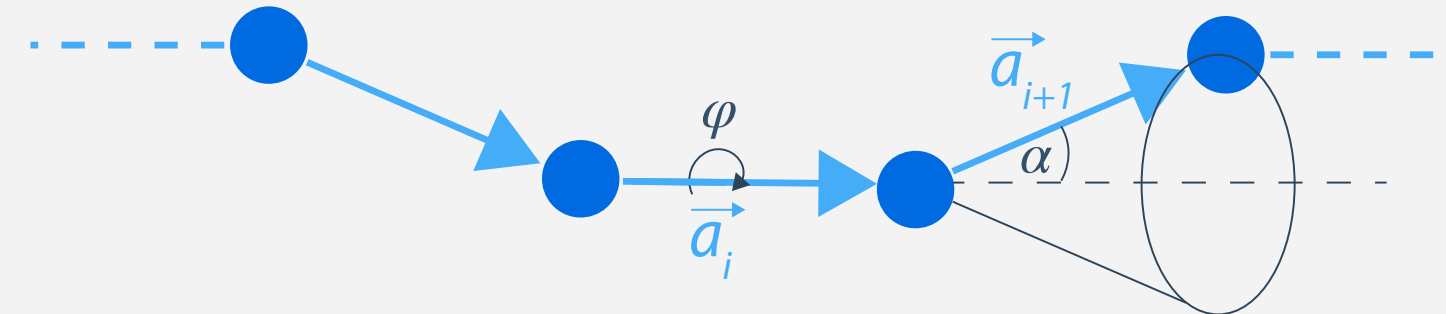
freely jointed chain



no torsion and bond angle restriction

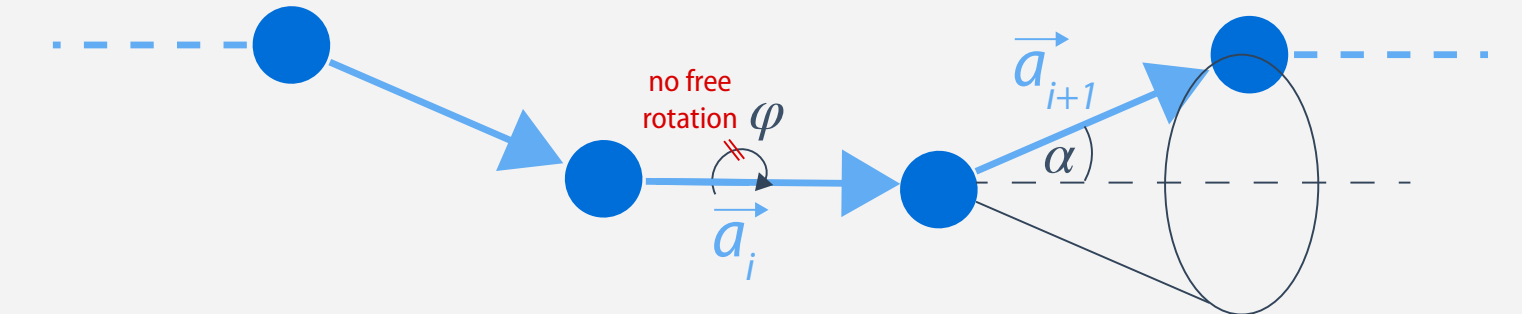
$$C_\infty = 1$$

freely rotating chain



bond angle restriction

hindered rotating chain



torsion and bond angle restriction

increased chain rigidity